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EPR study of Ti^{3+} ions formed under beta irradiation in silicate glasses

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Abstract: Reduction processes of Ti^{4+} ions in various silicate glasses have been studied by EPR spectroscopy at 20 K. Different parameters like the $[Na]/[Ti]$ ratio, the integrated dose and Ti^{3+} ions ageing processes were analyzed in this work. Simulation of the Ti^{3+} ion EPR spectra in different Ti-doped silicate glasses has shown three different Ti^{3+} environment attributed to one $^{[VI]}Ti^{3+}$ and two $^{[V]}Ti^{3+}$ environment (square pyramid and trigonal bi-pyramid). The $^{[VI]}Ti^{3+}$ ion environment is observed only for $[Na]/[Ti]$ ratios higher or equal to 68 although the two others are observed for all values of the $[Na]/[Ti]$ ratio considered. In terms of relative proportions, the Ti^{3+} ions in a square pyramid of oxygen are the dominant species.

1. Introduction

Reduction processes of Rare Earth (RE) and Transition Metals (TM) ions in silicate and borosilicate oxide glasses under ionizing irradiation have been observed by Electron Paramagnetic Resonance (EPR) and luminescence spectroscopies in previous works. Ions like Ti^{4+} [1], Sm^{3+} [2], Cr^{3+} [3] or Fe^{3+} [4] are reduced as a function of the integrated dose and complete reduction of Sm^{3+} ions has been

detected for integrated dose around 10^9 Gy using luminescence measurements [2]. The reduction processes of ions in glasses can be associated to the production of electron hole pairs (excitons) in glasses. Some electrons can be trapped for example on the Ti^{4+} ions, leading to the formation of stable Ti^{3+} ions while holes are trapped by oxygen ions leading to the production of different punctual holes defects in the irradiated glass [5]. Moreover, for the highest doses integrated inside the glass ($\approx 1 \times 10^9$ Gy), an increase of the network polymerization [6-11] associated with a production of molecular oxygen [6, 8-9, 11] may occur. These two processes are correlated to the alkaline migration processes under ionizing radiation [8, 10]. Efficiency of reduction processes under ionizing radiation should also depend on the environment around the doping ion in glasses.

Titanium ion is a good candidate for studying this relationship. First, when TiO_2 is added to a glass composition, the main oxidation state in the matrix is the Ti^{4+} state [12, 13]. Study of the titanium environment before irradiation will therefore bring information of the different environments around an oxidation state. The second reason is that titanium environment in glasses has been extensively studied by X-Ray Absorption Near Edge Structure (XANES), Extended X-ray Absorption Fine Structure (EXAFS) [12-18] and Raman spectroscopy [19, 20]. These works show the existence of different Ti^{4+} coordination number: $^{[IV]}Ti^{4+}$, $^{[V]}Ti^{4+}$ and $^{[VI]}Ti^{4+}$. The $^{[IV]}Ti^{4+}$ corresponds to a Ti^{4+} surrounded by a oxygen tetrahedron; the $^{[V]}Ti^{4+}$ to a Ti^{4+} ion in a square pyramid of oxygen with one short titanyl $Ti = O$ bond (1.7 Å) and four longer $Ti - O$ bonds (1.93 Å) [13]. The $^{[VI]}Ti^{4+}$ is associated to an octahedral environment around Ti^{4+} ions [12, 13]. The proportion between these environments is mainly correlated to the content of network modifiers inside the glass [13-14, 17-19, 21]. In addition, alkaline ions act as charge compensator to stabilize Ti^{4+} ions in both five-

and six-fold environment. Two reduction processes are therefore supposed to occur as a function of the integrated dose. The first one is a reduction process from Ti^{4+} to Ti^{3+} ions without any coordination number changes for doses lower than 1×10^9 Gy. The second one is associated to the alkaline migration at higher doses which could possibly lead to reduction processes of Ti^{4+} ions with a coordination number change. Under these assumptions and depending on the chemical composition of the glass, we shall observe in β -irradiated $SiO_2 - Na_2O - TiO_2$ glasses three different environments for Ti^{3+} ions: $[IV]Ti^{3+}$, $[V]Ti^{3+}$ and $[VI]Ti^{3+}$.

The aim of this paper is to study the different environments of Ti^{3+} ions produced under β -irradiation in $SiO_2 - Na_2O - TiO_2$ glasses with 1 to 4 %mol. of titanium as a function of the alkaline content and of the integrated dose. The second goal is to correlate these results on Ti^{3+} ions with literature results on Ti^{4+} environment as a function of the silicate glasses composition.

For that purpose, we have mainly used the X-band Electron Paramagnetic Resonance (EPR) spectroscopy at low temperature (20 K) associated with the simulation of experimental Ti^{3+} ions EPR spectra. This approach allows us to determine the EPR parameters (g -tensor, widths) for each Ti^{3+} ion sites produced under ionizing radiation and their relative proportions as the function of the glass compositions and integrated doses.

2. Experimental

2.1 Sample Preparation

The different glass compositions were prepared by mixing appropriate amounts of analytical reagent grade (SiO_2 , TiO_2 , Na_2CO_3) in an agate mortar with ethanol. The powders were then transferred in a Pt-crucible and placed in an air-oven at $850^\circ C$ for

10 h. This temperature step was used to ensure a correct decarbonation of the powders. After decarbonation, the melt was heated to 1600°C in 3 h and quenched. The obtained samples were cut and manually polished on a Labram HR Mecapol. The different glasses studied in this work are presented in table 1 where the sample number corresponds to the [Na]/[Ti] ratio of the silicate composition. The S10 ([Na]/[Ti] = 10), S20 ([Na]/[Ti] = 20) and S68 ([Na]/[Ti] = 68) samples are of the same composition than samples already studied by XANES and micro-Raman spectroscopy in previous works (resp. in [21] for S10 and S20 and [19] for S68).

All the samples were β -irradiated on a Van de Graaff accelerator (LSI, Ecole Polytechnique, France) with a 10 μ A current and an energy of 2.5 MeV at different doses: 3.24×10^6 ; 3.23×10^7 ; 2.75×10^8 ; 5.53×10^8 and 1.68×10^9 Gy. During irradiation, the sample temperature never exceeded 80°C.

All the samples were synthesized using the same protocol but the theoretical compositions has not been confirmed by a chemical analysis. Thus, we'll assume that the composition of all samples is affected by the same error.

2.2. Experimental EPR spectroscopy

X-band ($\nu = 9.490$ GHz) EPR spectra were recorded at 20 K using an OXFORD He cryostat on a EMX Bruker spectrometer. The microwave power used, 0.05 mW, were determined by considering the saturation properties of Ti^{3+} EPR lines at 20 K. The different EPR spectra presented in this paper have been normalized to the same receiver gain (1.10^4) and sample weight (100 mg) in order to get quantitative comparison of the EPR lines intensities between the different samples. We emphasize that EPR spectroscopy is a tool where the absorbance spectra (i.e. the experimental spectrum integrated twice) of all EPR lines corresponding to one

paramagnetic ion in a specific site can be correlated to the paramagnetic specie spin content inside the sample [22]. However, due to different saturation properties, the EPR linewidths of known EPR reference spectra (e.g. DPPH, strong and weak pitch) and Ti^{3+} EPR spectra are strongly different. In the present work it is thus very difficult to compare these spectra to determine the absolute number of Ti^{3+} spins. Then, in this study, we will only use arbitrary units (a.u.) to analyze the evolution of Ti^{3+} ions content as a function of the integrated dose and silicate glass compositions. However, using simulation of EPR spectra (see §2.3), we are able to separate the experimental spectrum into a sum of Ti^{3+} elementary spectra corresponding to different sites of Ti^{3+} ions. In this case, the relative proportion of Ti^{3+} in each site, expressed in %, will be determined by computing the ratio of the area under each elementary absorbance spectrum and the area under the total simulated spectrum.

Many different phenomena may induce an error on the measured values: position of the sample into the EPR cavity, variation of temperature, etc. The global error on each point of the EPR spectrum is estimated to be 4.8 %. This value was obtained by recording ten times the EPR spectrum of the same sample (S80 irradiated at $3,27 \times 10^7$ Gy). Each time the sample is removed and replaced into the EPR cavity and the temperature returns to ambient before decreased again.

2.3. Simulation of EPR spectra

Simulation of EPR spectra were realized using the ZFSFIT program developed by G. Morin and D. Bonnin [23]. This program was specially developed for fitting EPR powder spectra with paramagnetic species with nuclear spin $I \leq 7/2$ and electronic spin $S \leq 5/2$. The more general spin Hamiltonian calculated by ZFSFIT is:

$$\beta H \cdot g \cdot S + \sum_{k=2,4} \sum_{q=-k}^{+k} B_k^q O_k^q + S \cdot A \cdot I$$

Where β is the Bohr magnetron, H is the magnetic field tensor, g is the Lande tensor, S the spin tensor, $O_k^{\mp q}$ the Stevens operators, $B_k^{\mp q}$ the constant associated with the Stevens operators, A the hyperfine tensor and I the nuclear spin operator [23]. For trivalent titanium ions, $S = 1/2$ and $I = 0$. Therefore, the spin Hamiltonian used in our simulations is:

$$\beta H \cdot g \cdot S$$

Because we observe different Ti^{3+} EPR sites in our glasses, the simulated spectra calculated by ZFSFIT is the sum of the elementary spectrum of Ti^{3+} ions in each of these sites. ZFSFIT compares the calculated spectrum with the experimental one, compute the difference between both and refine the given parameters (g values, isotropic width (s_0), and anisotropic widths (s_1 , s_2 and s_3) for each Ti^{3+} ion EPR site) by a routine based on the Levenberg-Marquardt algorithm.

At the end of the fitting procedure, the obtained simulated EPR spectra for each Ti^{3+} environment are defined by three g -values which can be associated to the EPR axis of the system [24]. Because we are working with amorphous materials, our EPR spectra are the sum over all spatial orientations of the Ti^{3+} EPR spectrum and thus, it is impossible to associate a g value with a spatial direction into the sample. By convention, we will assume that $g_3 \geq g_2 \geq g_1$. The three anisotropic widths s_1 , s_2 and s_3 widths correspond to the g -strain process in amorphous material associated with g values distribution. The g -strain process in the ZFSFIT program is treated as a perturbation of the Zeeman Hamiltonian term different along the g tensor direction. By convention, we will assume $s_3 = 0$ to precisely quantify the other widths.

3. Results

3.1. Influence of $[Na]/[Ti]$ ratio

The figure 1 presents the EPR spectra of titanium-doped silicate glasses samples S10 to S80 after β -irradiation at 3.23×10^7 Gy. These spectra are complex and may be separated into two groups of lines. The first group, centered around $g = 2.0076 \pm 0.003$ and close to the g -value of free electron ($g_e = 2.0023$ [22]) is attributed to the paramagnetic hole defects created by irradiation [7, 5] and the second group, centered around $g = 1.935 \pm 0.003$ is attributed to Ti^{3+} ions [22, 25, 26]. In this paper we will mainly focus our study on the EPR spectrum of Ti^{3+} ions. Fig. 1 clearly shows a dependence between the shape of the Ti^{3+} EPR spectra and the $[Na]/[Ti]$ ratio. For a ratio less than 40, all the EPR spectra are identical and are characterized by a positive component at $g = 1.935 \pm 0.003$ and a large negative component around $g \approx 1.855$. When the $[Na]/[Ti]$ ratio increases to 80, the EPR spectra shows an additional line determined by an intense positive peak at $g = 1.963 \pm 0.003$ and a negative shoulder at $g = 1.935 \pm 0.003$. Considering the existence of different EPR lines on the Ti^{3+} EPR spectra, we can conclude first that the Ti^{3+} EPR signal of these irradiated silicate glasses is the sum of different components corresponding to, at least, two different Ti^{3+} EPR sites. The first Ti^{3+} site in silicate glasses, hereafter denoted s_I component in the EPR spectra, appears for $[Na]/[Ti]$ ratios greater or equal to 68 and can be determined by $g_{//} \approx 1.963$ and $g_{\perp} \approx 1.855$ axial symmetry. The second component in the Ti^{3+} EPR spectra is observed for all $[Na]/[Ti]$ ratios. However, we will show later that this EPR line is in fact the sum of two different components denoted s_{II} and s_{III} corresponding to two other Ti^{3+} environments in the irradiated silicate glasses.

3.2. Influence of total integrated dose

Figure 2 shows the EPR spectra of the two titanium doped silicate glasses S20 and S80 irradiated at different doses between 3.24×10^6 and 1.68×10^9 Gy. The EPR spectra of pristine S20 and S80 samples also presented in Fig. 2 don't show any signal in the region of Ti^{3+} EPR lines. We can conclude therefore that Ti^{3+} ions for these silicate glass compositions are only produced during irradiation and not during the glass synthesis. For the S20 sample ($[Na]/[Ti]=20$), we observe an increase of the Ti^{3+} EPR lines intensity corresponding to an increase of the Ti^{3+} content as a function of the dose integrated in the glass. We can also analyze the appearance of a shoulder at $g = 1.963 \pm 0.003$ and a shift of the minimum of the EPR signal from $g = 1.846 \pm 0.003$ to $g = 1.871 \pm 0.003$. According to our previous results, we can correlate these changes in the Ti^{3+} EPR spectra to a modification of the relative proportion of Ti^{3+} ions between the s_{II} and s_{III} sites when the integrated dose increases. We can then attribute the s_{II} component EPR spectrum to the shoulder at $g = 1.963$ and the minimum at $g = 1.871$. Finally, the last Ti^{3+} environment characterized by the s_{III} component has an EPR spectrum close to the EPR spectra observed for the lowest dose (3.24×10^6 Gy).

For the S80 sample, we observe also an increase of the Ti^{3+} EPR spectrum intensity with the integrated dose inside the silicate glass correlated with the increase of Ti^{3+} ions content. We can also see on Fig. 2 an increase of the relative intensity of $g = 1.935 \pm 0.003$ component and a displacement of the minimum of the signal from $g = 1.855 \pm 0.003$ to $g = 1.868 \pm 0.003$ for integrated dose in silicate glasses higher than 5.58×10^8 Gy. This result can be correlated to a decrease of the number of Ti^{3+} ions in the s_I site ($g_{//} \approx 1.963$ and $g_{\perp} \approx 1.855$) when the integrated dose is in the order of the giga-gray. At these higher integrated doses, migration of alkaline ions under ionizing radiation has been observed [8, 10]. The decrease of the s_I component could

be correlated to the migration of alkaline ion acting as charge compensator of the Ti^{3+} ions in this environment. Thus, this result could imply that the coordination number of Ti^{3+} ion in s_I site is greater than the coordination number of both Ti^{3+} in s_{II} and s_{III} sites.

3.3. *Post irradiation ageing processes.*

The figure 3 presents the EPR spectra of the S20 and S80 samples recorded just after their β -irradiation at 3.23×10^7 Gy and 165 days later. No significant changes are analyzed on the EPR spectra of the S20 sample as a function of the ageing time. By contrary for the S80 sample, we observe an important modification of the Ti^{3+} EPR signal mainly correlated to the intensity decrease of the s_I component. The glass sample ageing is also associated to a visible leaching of the S80 surface sample. Glass corrosion by water leads to three different kinds of process [27, 28]: hydration, hydrolysis and ionic exchange. In particular, ionic exchange between Na^+ and H^+ ions during ageing processes lead to a decrease of the charge compensation intensity around Ti ions and thus, to a reorganization of Ti^{3+} environment. In the present case, we interpret this reorganization as a decrease of the coordination number around Ti^{3+} . The similarities between the evolutions of the EPR spectrum of Ti doped silicate glass irradiated at high-dose (close or higher to 1×10^9 Gy) showed on figure 2 and after a 165 days ageing presented on figure 3 confirm that the s_I site of Ti^{3+} ions has a higher coordination number than the other Ti^{3+} sites in the irradiated silicate glasses studied in this work.

3.4. *Simulation of experimental Ti^{3+} EPR spectra*

Using our previous observations on figures 1-3, we tried to reproduce the experimental EPR spectra of all our samples β -irradiated at different doses between 3.24×10^6 and 1.68×10^9 Gy. To determine the EPR parameters of the s_I site, we subtract the EPR spectra obtained 165 days after irradiation to the EPR spectra recorded just after irradiation. For the other sites, we tried different assumptions and we only kept the solution which reproduce the better all our EPR spectra. We conclude to the existence of three different EPR sites of Ti^{3+} determined by the EPR parameters given in table 2. Figure 4 shows for example the experimental X-band EPR spectrum, the three s_I , s_{II} and s_{III} components used to obtain the simulated spectrum and the total simulated spectrum defined by the sum of the s_I , s_{II} and s_{III} components for the S80 sample.

The obtained s_I , s_{II} and s_{III} EPR spectra are very close to some results already reported in previous studies. For example, S. Arafa and F. Assaghy [26] observed in X-irradiated $SiO_2 - Na_2O - TiO_2$ glasses an EPR spectrum characterized by $g_{\perp} = 1.9693$ and $g_{\parallel} = 1.9413$ very similar to our s_I spectrum. H. Böhm and G. Bayer [29] also obtained a similar spectrum defined by $g_{\perp} = 1.970$ and $g_{\parallel} = 1.930$ by studying the $[^{VI}]Ti^{3+}$ formed under reducing atmosphere in Mg_2TiO_4 crystals. V. Laguta *et al.* [30] reported the EPR spectrum of Ti^{3+} ions surrounded by five oxygen atoms and one oxygen vacancy in $PbTiO_3$ crystals. Their spectrum, characterized by $g_1 = 1.965$; $g_2 = 1.919$ and $g_3 = 1.848$ is very close to our s_{II} spectrum. And H. Böhm and G. Bayer [29], also published the EPR spectrum of Ti^{3+} in La_2TiO_5 crystals which is close to our s_{III} spectrum. All these previous studies supports the simulation solution we've kept. However the shape of the simulated EPR spectra will be discussed in section 4.

3.5. Proportions of Ti^{3+} ions in each site

By a double integration of the simulated EPR spectra, we can obtain the proportion of Ti^{3+} ions in each titanium site inside the glass. Figure 5 presents the evolution of the mean proportion over all samples of Ti^{3+} ions in each site as a function of the total integrated dose (in log. scale). According to this figure, the Ti^{3+} ions produced during ionizing radiation in the silicate glasses are mainly in the s_{III} site followed by the s_{II} site and finally the s_I site which is only observed in the S68 and S80 samples. Moreover, Ti^{3+} ions content in the s_I site seems to be constant or, as we observed in figure 2, decreases when the dose is close or higher than the giga-gray. As we already said, some of our samples have the same chemical composition as samples already studied in previous works. First of all, N. Ollier *et al.* [1] showed by XANES spectroscopy the existence of a mixture of $^{IV}Ti^{4+}$ and $^{VI}Ti^{4+}$ in a sample very close to our S68 sample but with a $[Na]/[Ti]$ ratio equals to 66. More precisely, in his XANES spectroscopy study [21], F. Farges concludes to the existence in the S10 sample of 80 % of $^{IV}Ti^{4+}$ and 20 % of $^{VI}Ti^{4+}$. Depending on the total integrated dose, we observe in the S10 sample from 20 to 30 % of Ti^{3+} ions in the s_{II} site and from 70 to 80 % of Ti^{3+} ions in the s_{III} site. In the case of the S20 sample, F. Farges concluded to the existence of 85% of $^{IV}Ti^{4+}$ and 15% of $^{VI}Ti^{4+}$. In figure 5, we observe from 80 to 90% of Ti^{3+} ions in the s_{III} site and from 10 to 20% of Ti^{3+} ions in the s_{II} site. With the assumption of a direct reduction from Ti^{4+} to Ti^{3+} for doses lower than 1×10^9 Gy, we must emphasize the similarities between 1) the quantities of $^{IV}Ti^{4+}$ and the quantities of Ti^{3+} ions in the s_{III} site and 2) the quantities of $^{VI}Ti^{4+}$ and the quantities of Ti^{3+} ions in the s_{II} site. Then, we conclude that the Ti^{3+} ions in the s_{III} site comes from the reduction of $^{IV}Ti^{4+}$ ions and the Ti^{3+} ions in the s_{II} comes from the reduction of $^{VI}Ti^{4+}$ ions. The question of the reduction efficiency will be discussed in the next section.

4. Discussion

Previous works in literature have mainly concluded to the presence in glasses of Ti^{3+} ion in an octahedral environment even if authors observe large differences in g tensor values derived from the studies. For example, Y. Kim and P. Bray analyzed in γ -irradiated $K_2O - TiO_2$ glasses [31] an $^{VI}Ti^{3+}$ EPR spectrum characterized by $g_{\perp} = 1.975$ and $g_{\parallel} = 1.890$. S. Arafa and A. Bishay [32] attributed also to $^{VI}Ti^{3+}$ a very symmetric Ti^{3+} EPR spectrum characterized by $g \approx 1.942$ in γ -irradiated $B_2O_3 - CaO - TiO_2$ glasses. However, our EPR results suggest the existence of different Ti^{3+} sites relative to different coordination numbers for Ti^{3+} ions. As shown in § 3.2, the modifications of the EPR spectrum for the samples irradiated at the highest doses suggest a relation between the existence of Ti^{3+} ions in the s_I site and the number of sodium ions available for charge compensation. This is confirmed by the fact that we observe this site only on the EPR spectra of the samples with the highest $[Na]/[Ti]$ ratio (e.g. S68 and S80) and by the lower stability of this site during leaching (§ 3.4). We concluded that the coordination of the titanium ions in the s_I site is higher than the coordination of Ti^{3+} ions in the s_{II} and s_{III} sites. To our knowledge, the higher observed coordination number for Ti^{3+} ions is VI. Thus, we attribute the s_I site to $^{VI}Ti^{3+}$ ions.

F. Farges [21] has proposed a possible square based pyramidal environment ($^{IV}Ti^{4+}$) in silicate glasses. In this coordination number [13, 16, 21], the $^{IV}Ti^{4+}$ ions environment is characterized by one $Ti = O$ titanyl bond and four $Ti - O$ bonds. The assumption of a direct reduction processes for Ti doped silicate glasses studied in this work suppose production of $^{VI}Ti^{3+}$ ions in this square pyramidal environment. This particular stronger titanyl bond will influence strongly the powder EPR spectrum

linewidth of this Ti^{3+} site. The width of Ti^{3+} EPR line may be understood as the result of a *g-strain* process (i.e. a distribution of the *g*-values) [24]. This distribution traduces the variation of the environment of the paramagnetic ion along specific directions of the *g* tensor. Thus, we could therefore correlate the particular shape of the $s_{///}$ EPR spectrum (Table 2, Fig. 4) characterized by a small width (0.00515 cm^{-1}) for the g_z component of the Ti^{3+} powder spectrum in comparison to the g_y (0.06273 cm^{-1}) and g_x (0.13127 cm^{-1}) width values to $^{VI}\text{Ti}^{3+}$ ions in square base pyramidal environment. Moreover, $^{VI}\text{Ti}^{3+}$ ions were already observed in some previous studies, for example by M. Nolan *et al.* [33] in non-stoichiometric TiO_2 rutile.

The site $s_{//}$ is more difficult to associate with a coordination number around Ti^{3+} ions. As we said it, we concluded that the coordination number associated to $s_{/}$ is higher than the coordination number associated to the $s_{//}$ site. However, we observed by studying the proportions of Ti^{3+} in each site (§ 3.5) that $s_{//}$ may correspond to a coordination number of VI. To explain this apparent discrepancy of our results, we propose to associate the $s_{//}$ site of Ti^{3+} with a coordination number of V in a trigonal bi-pyramid shape (i.e. an octahedron with an oxygen vacancy). Such geometry was for example already observed for Ti^{4+} in La_2TiO_5 crystals and it is consistent with the EPR results of V. Laguta *et al.* [30]. F. Farges [20] didn't take into account this geometry in his XANES study because it is less common to observe Ti^{4+} in trigonal bi-pyramide than in square pyramid. Moreover, his results are defined with a $\pm 10 \%$ error. Thus, we can assume that the trigonal bi pyramid geometry is an intermediary case between the octahedron and the square pyramid which was not taken into account in the previous studies. This may explain the apparent discrepancy between our results. However, because figure 5 shows different slopes of reduction, it may exist in our samples some different reduction efficiency, depending on the titanium

sites and on the integrated dose. This different reduction efficiency may change our result, especially concerning the total Ti^{3+} ions quantities. This problem must be taken into account for our next studies.

5. Conclusion

We studied in this work by X-band EPR spectroscopy and simulation of EPR spectra the environment of Ti^{3+} ions produced under β -irradiation in Ti-doped $\text{SiO}_2 - \text{Na}_2\text{O} - \text{TiO}_2$ glasses. We conclude to the existence of three different Ti^{3+} EPR sites, denoted s_I , s_{II} and s_{III} . The first one, s_I ($g_{//} = 1.970$ and $g_{\perp} = 1.939$), is associated to a $^{VI}\text{Ti}^{3+}$ in an octahedron of oxygens and is observed only when the $[\text{Na}]/[\text{Ti}]$ ratio of the glass is higher or equal to 68 and the integrated dose less than 1×10^9 Gy. The second site, s_{II} ($g_1 = 1.868$, $g_2 = 1.904$, $g_3 = 1.965$) is attributed to a Ti^{3+} ion in a trigonal bi-pyramid environment. The last one, s_{III} ($g_1 = 1.763$, $g_2 = 1.863$, $g_3 = 1.936$), is dominant (80 % of the total amount of observed Ti^{3+} ions) and corresponds to a $^{VI}\text{Ti}^{3+}$ ion in a square pyramid of oxygens. The s_{II} and s_{III} sites are observed for all values of the $[\text{Na}]/[\text{Ti}]$ ratio.

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7. References

- [1] N. Ollier, P. Lombard, F. Farges, B. Boizot, *Journal of Non Crystalline Solids*, **354**, 480 (2008).
- [2] E. Malchukova, B. Boizot, G. Petite, D. Ghaleb, *Journal of Non-Crystalline Solids* **353**, 2397 (2007).
- [3] B. Boizot, F.Y. Olivier, G. Petite, D. Ghaleb, *Nuclear Instrument and Methods in Physics Research B* **266**, 2966 (2008).

- [4] F.Y. Olivier, B. Boizot, D. Ghaleb, G. Petite, *Journal of Non-Crystalline Solids* **351**, 1061 (2005).
- [5] B. Boizot, G. Petite, D. Ghaleb and G. Calas. *Nuclear Instrument and Methods in Physics Research B* **141**, 580 (1998).
- [6] N. Ollier, B. Champagnon, B. Boizot, Y. Guyot, G. Panczer and B. Padlyak, *Journal of Non Crystalline Solids* **323**, 200 (2003).
- [7] B. Boizot, N. Ollier, F. Olivier, G. Petite, D. Ghaleb and E. Malchukova. *Nuclear Instrument and Methods in Physics Research B* **240**, 146 (2005).
- [8] B. Boizot, G. Petite, D. Ghaleb, B. Reynard and G. Calas, *Journal of Non Crystalline Solids* **243**, 268 (1999).
- [9] N. Ollier, B. Boizot, R. Reynard, D. Ghaleb and G. Petite, *Journal of Nuclear Materials* **340**, 209 (2005).
- [10] B. Boizot, G. Petite, D. Ghaleb, N. Pellerin, F. Fayon, B. Reynard and G. Calas, *Nuclear Instruments and Methods in Physics Research B* **166**, 500 (2000).
- [11] D. Neuville, L. Cormier, B. Boizot, A. Flank, *Journal of Non Crystalline Solids*, **323**, 207 (2003).
- [12] R. Gregor, F. Lytle, D. Sandstrom, J. Wong and P. Schultz, *Journal of Non Crystalline Solids* **55**, 27 (1983).
- [13] F. Farges, G. Brown, E. Navrotsky, H. Gan and J. Rehr, *Geochimica et Cosmochimica Acta* **60**, 3039 (1996).
- [14] G. Henderson and J. St – Amour, *Chemical Geology* **213**, 31 (2004).
- [15] G. Henderson and M. Fleet, *Journal of Non Crystalline Solids* **211**, 214 (1997).
- [16] F. Farges, G. Brown and J. Rehr, *Geochimica and Cosmochimica Acta* **60**, 3023 (1996).
- [17] F. Farges, *Journal of Non Crystalline Solids* **244**, 25 (1999).
- [18] G. Henderson, X. Liu and M. Fleet, *Mineralogical Magazine* **67**, 597 (2003).
- [19] B. Mysen and D. Neuville *Geochimica and Cosmochimica Acta* **59**(2), 325 (1995).
- [20] G. Henderson, *The Canadian Mineralogist* **33**, 399 (1995).
- [21] F. Farges. *American Mineralogist* **82**, 36 (1997).
- [22] A. Abragam and B. Bleaney. *Electron Paramagnetic Resonance of transition ions*. Clarendon Press, 1970.
- [23] G. Morin and D. Bonnin. *Journal of Magnetic Resonance* **136**, 176 (1999).
- [24] A. Weil, J. Bolton and J. Wertz. *Electron Paramagnetic Resonance*. Wiley Interscience, 1994.

- [25] N. Iwamoto, H. Hidaka and Y. Makino. *Journal of Non-Crystalline Solids* **58**, 131 (1983).
- [26] S. Arafa and F. Assabghy. *Journal of Applied Physics* **45**(12) 5269 (1974).
- [27] C. Cailleteau, C. Weigel, A. Ledieu, P. Barboux and F. Devreux. *Journal of Non-Crystalline solids* **354**, 117 (2008).
- [28] B. Bunker. *Journal of Non-Crystalline Solids* **179**, 300 (1994).
- [29] H. Bohm and G. Bayer, *Journal of Non Crystalline Solids*, **31**, 2125 (1970). [30] V. Laguta, M. Glinchuk, I. Bykov and Y. Maksimenko, *Physical Review B* **54**, 12353 (1996).
- [31] Y. Kim and P. Bray, *The Journal of Chemical Physics* **53**, 716 (1970).
- [32] S. Arafa and A. Bishay, *Physics and Chemistry of Glasses* **11**, 75 (1970).
- [33] M. Nolan, S. Elliott, J. Mulley, R. Bennett, M. Basham and P. Mulheran, *Physical Review B* **77**, 235424 (2008).

Figure captions

Figure 1. X-band EPR spectra recorded at 20 K of $\text{SiO}_2 - \text{Na}_2\text{O} - \text{TiO}_2$ glasses β -irradiated at 3.23×10^7 Gy. The g values indicated in the figures correspond to the minima or maxima of the different Ti^{3+} EPR lines.

Figure 2. X-band EPR spectra recorded at 20 K of the S20 and S80 samples non-irradiated and β -irradiated at different doses between 3.24×10^6 and 1.68×10^9 Gy. The g values indicated in the figures correspond to the minima or maxima of the different Ti^{3+} EPR lines.

Figure 3. X-band EPR spectra at 20 K of the S20 and S80 samples recorded just after (“initial”) and 165 after β -irradiation at 3.23×10^7 Gy. The g values indicated in the figures correspond to the minima or maxima of the different Ti^{3+} EPR lines.

Figure 4. Experimental, simulated and components of the X-band EPR spectrum of the S80 sample β -irradiated to 3.24×10^7 Gy.

Figure 5. Mean proportion of the number of Ti^{3+} ions in the s_I , s_{II} and s_{III} sites identified in β -irradiated titanium doped silicate glasses as the function of the dose integrated.

Tables

Sample	[SiO ₂]	[Na ₂ O]	[TiO ₂]	[Na]/[Ti]	Ref.
S10	76	20	4	10	[21]
S20	78	20	2	20	[21]
S40	79	20	1	40	-
S68	65	34	1	68	[19]
S80	59	40	1	80	-

Table 1. Samples composition. All quantities are in mol%.

Site	g_1	g_2	g_3	s_0	s_1	s_2
s_I	1.939	1.970		0.00282	0	0
s_{II}	1.868	1.904	1.965	0.00363	0.03107	0.02088
s_{III}	1.763	1.863	1.936	0.00514	0.13127	0.06273

Table 2. EPR parameters determined by simulation for the three Ti³⁺ sites observed in β -irradiated SiO₂ – Na₂O – TiO₂ glasses. The s_0 , s_1 and s_2 width are in cm⁻¹.

Figure 1

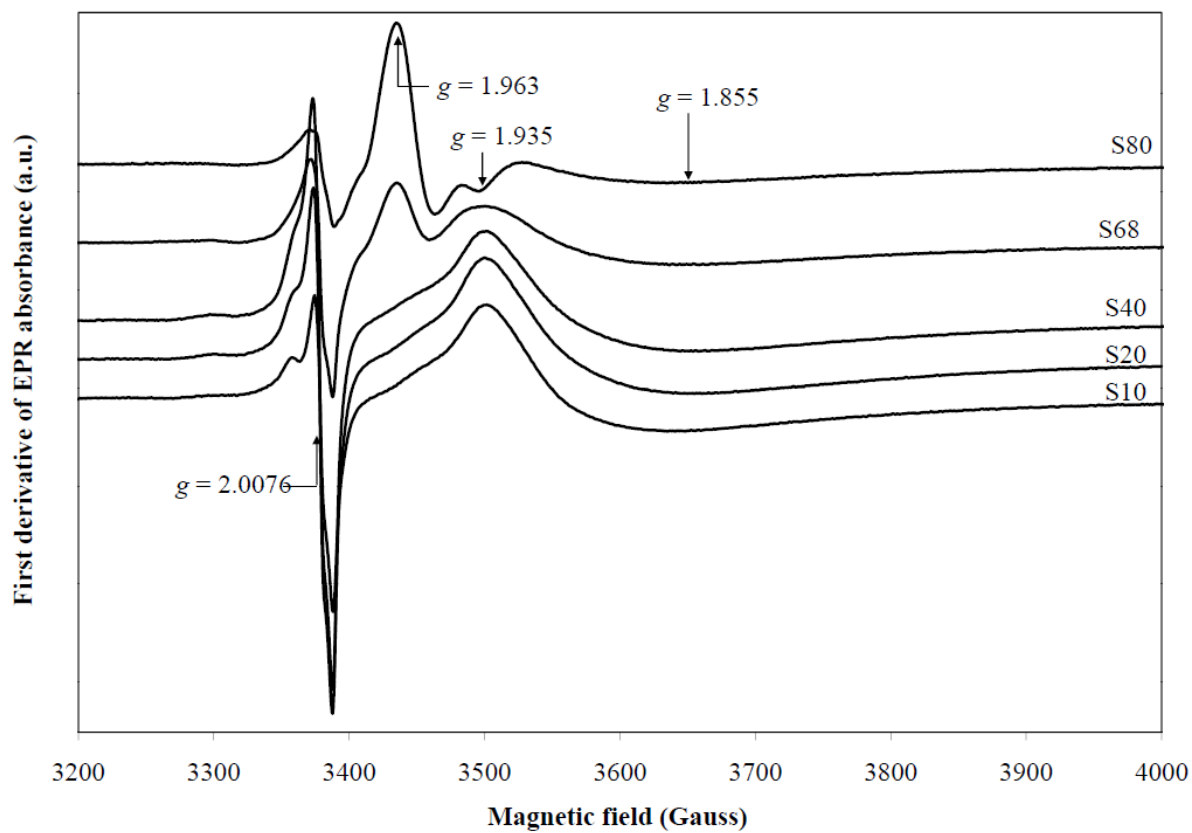


Figure 2

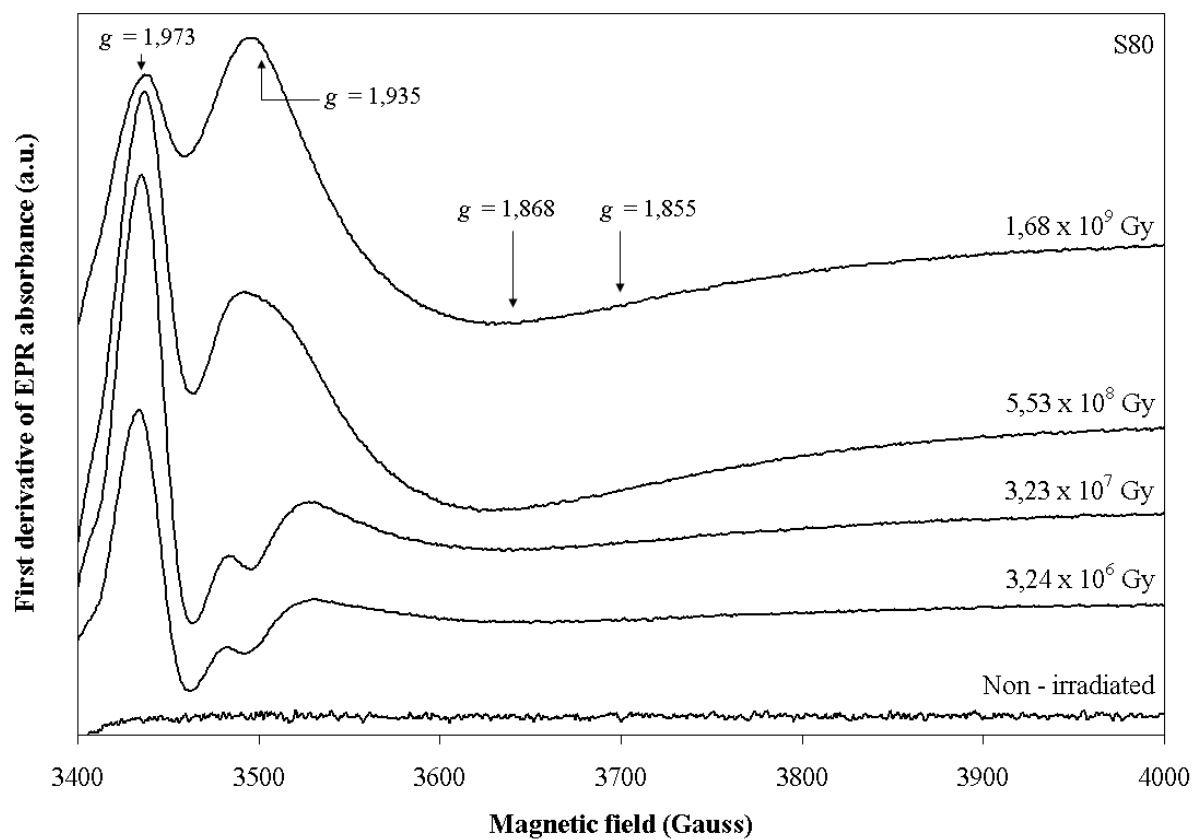
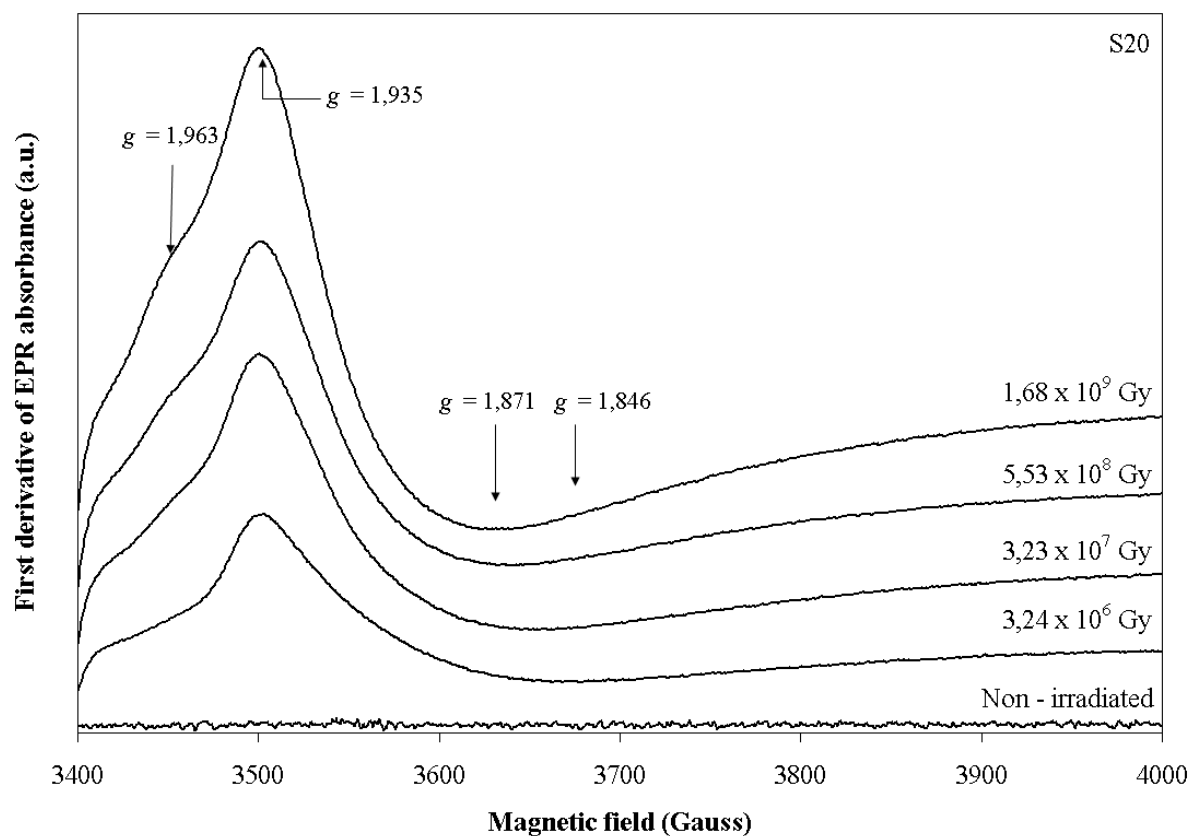


Figure 3

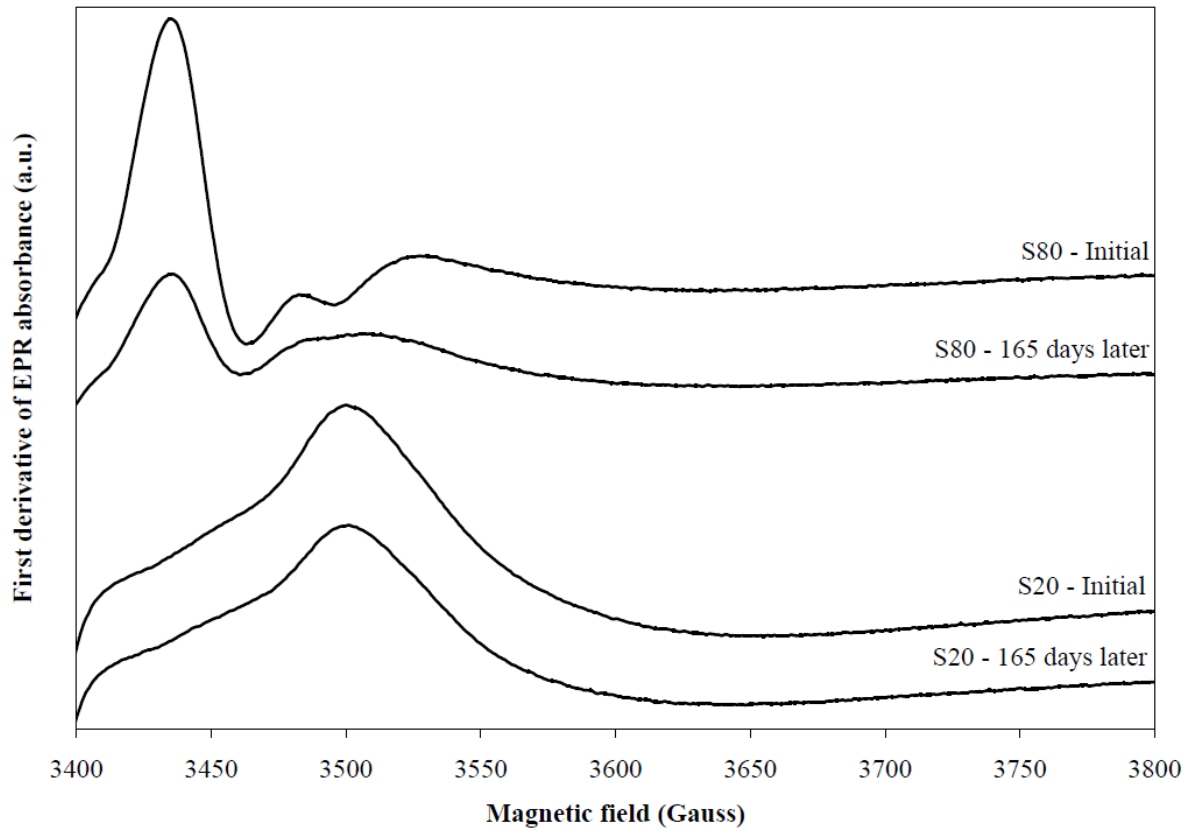


Figure 4

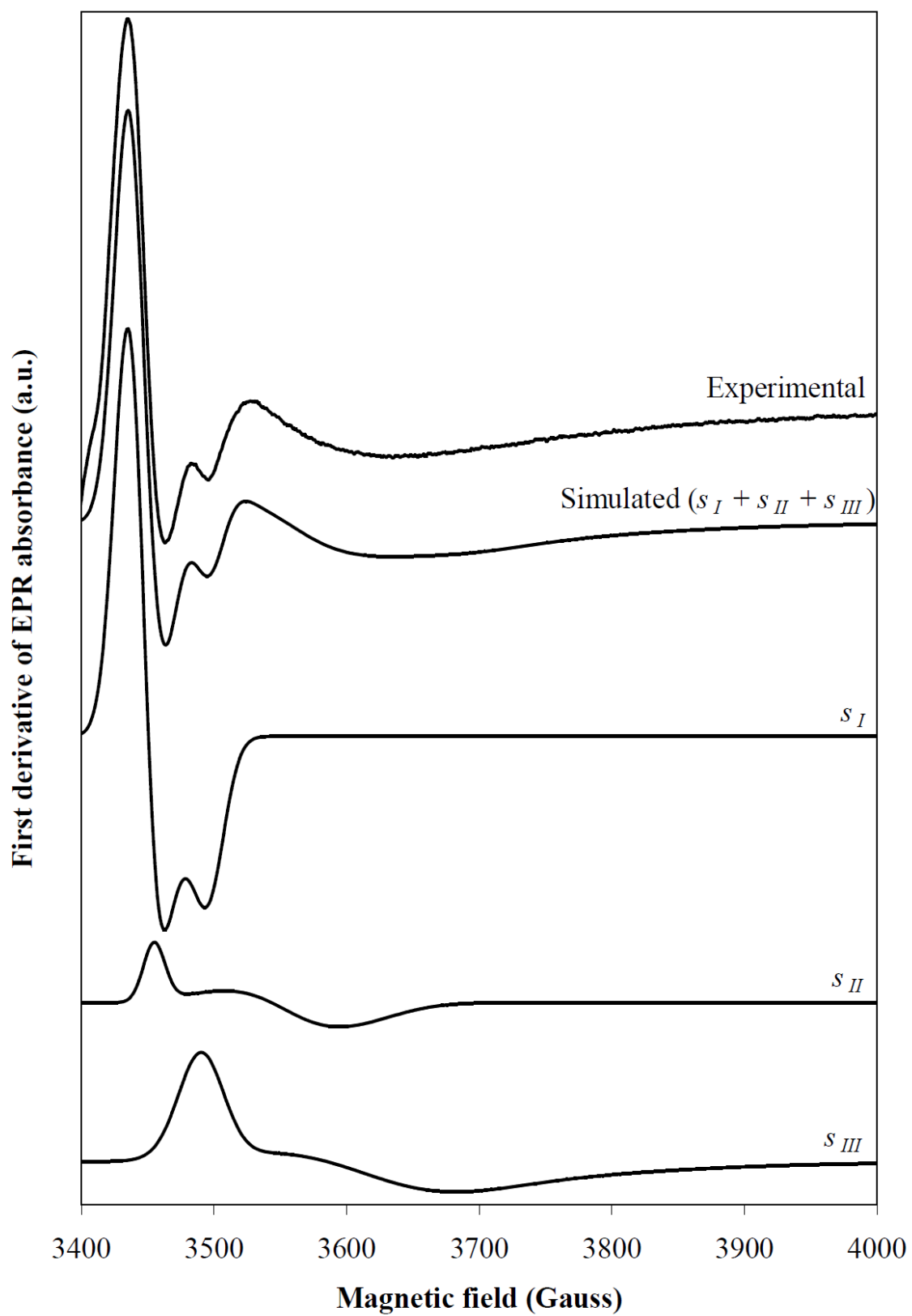


Figure 5

