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# Ciliary white-light: optical aspect of ultrashort laser ablation on transparent dielectrics

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## Abstract

We report on a novel nonlinear optical phenomenon, coined as ciliary white-light, during laser ablation of transparent dielectrics. It is observed in 14 different transparent materials including glasses, crystals, and polymers. This phenomenon is also universal with respect to laser polarization, pulse duration and focusing geometry. We interpret its formation in terms of the nonlinear diffraction of the laser generated white-light by the ablation crater covered by nanostructures. It carries rich information on the damage profile and morphology dynamics of the ablated surface, providing a real time *in-situ* observation of the laser ablation process.

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Laser ablation on dielectrics has attracted many interests since the early days of the laser [1-3]. With femtosecond laser pulses, heat diffusion to the surrounding of the irradiated area is greatly suppressed, because the pulse duration is shorter than the electron-phonon coupling time in the materials [2, 3]. As a consequence, femtosecond laser ablation induces minimal thermal side effect. On the other hand, femtosecond laser induced modifications can even be confined to sub-wavelength scale due to the nonlinear nature of energy deposition [3]. These unique features, combined with laser beam scanning or sample translation, have enabled a wide spectrum of applications, such as deep hole drilling on hard or brittle glass [3, 4], fabrication of versatile microfluidic devices [5], femtosecond laser in situ keratomileusis (LASIK) and dental surgery [6, 7], synthesis of nanomaterials [8], thin film deposition [2, 9], artwork preservation and restoration [10, 11], *etc.* The physical mechanism underlying ablation on dielectrics is quite different from that of metals and semiconductors. Dielectrics are often transparent to laser emission. Therefore, initial free electrons are usually generated via strong-field ionization, followed by impact ionization while the electrons are heated by the laser field through inverse Bremsstrahlung. In the presence of free electrons, the remaining laser radiation starts to be absorbed strongly, leading to further heating of the electron in which case dielectrics start to behave like metals. Subsequent electron diffusion, electron-lattice coupling, vapor expansion, surface macroscopic Coulomb explosion, add further complexities to this multifaceted process [2, 3, 9, 12, 13]. Even if a lot of studies have focused on the fundamental aspect of laser ablation [2, 12, 13], surprisingly little attention has been given so far to the evolution of the laser pulse itself during this surface ablation process.

Here we present the first experimental investigation on the laser pulses evolution in the spatial and spectrum domain during laser ablation of transparent materials. Radiation patterns in the form of an ellipse and a ring with large opening angle ( $\sim 100^\circ$ ) are observed, in addition to the directly transmitted laser beam. The elliptic radiation shows a laser polarization-dependent orientation while the ring-shaped does not. More interestingly, the ring-shaped radiation consists of numerous colored needle-like sub-structures with radial orientation. Owing to its spatial and spectral properties, we name this ring-shaped radiation “ciliary white-light” (CWL). We demonstrate the universality of the CWL with respect to the nature of materials (14 transparent media) and laser parameters (polarization states, pulse duration, and focusing geometry). This phenomenon, carrying abundant structural information of the ablation site, provides a real time *in-situ* observation of laser ablation on transparent materials.

Figure 1(a) presents a schematic of the experiments. A chirped pulse amplification (CPA) system (Thales, Alpha 100) is used as the main laser source. This system delivers 50 fs laser pulses with a maximum pulse energy of 10 mJ at a repetition rate of 100 Hz. Femtosecond laser pulses are focused by different convex focal lenses on the front surface of samples in ambient air. With focal lens of  $f = 1000$  mm, the beam waist is estimated to be 91  $\mu\text{m}$ . The polarization state of the laser pulses can be changed with a  $\lambda/2$  or a  $\lambda/4$  wave plate before the lens. The sample is

mounted on a mechanical 3-dimensional translation stage to insure a fresh site after each series of laser shots. A quasi-transparent tracing paper serving as a diffusing screen is installed 23 mm after the sample incident surface. A charge coupled device (CCD) camera captures the laser emission pattern on the screen.

For horizontally polarized laser pulses with energy of  $E_{in} = 15 \mu\text{J}$ , a stable, on axis, white light spot is observed on the screen even up to 200,000 laser shots (Fig. 1(b)). This white-light, also referred as supercontinuum emission, has been exploited in many applications such as laser spectroscopy and optical parametric chirped pulse amplification [14, 15]. Interestingly, beyond a threshold laser energy of  $E_{in} = 55 \mu\text{J}$ , a new phenomenon becomes apparent with increasing laser shot numbers  $N$ . Figure 2 presents the result for  $E_{in} = 210 \mu\text{J}$ . After  $N = 60$ , the stable white-light spot breaks up and a large number of flickering, colorful speckles appear around the central spot. With the further accumulation of laser shots, these colorful speckles self-organize into two structures: an inner ellipse and an outer ring (Fig. 2(c)-(f)). After 175 laser shots, the ellipse gradually disappears while the outer ring dominates (Fig. 2(e)-(i)). The ring-shaped radiation becomes quite stable after about 410 shots. It exhibits a divergence angle of about  $100^\circ$ , 9 times larger than the central white-light conical emission (Fig. 2(i)). In this stable state, the central white-light contains only 18.9% (18  $\mu\text{J}$ ) of the total transmitted laser energy (95  $\mu\text{J}$ ), indicating that 81.1% of laser energy is now contained in the CWL.

Several features are worth noticing in the above process. First, both the outer ring and inner ellipse structure become larger upon increase of laser shots. Fig. 3(a) shows their opening angles as a function of  $N$ . Second, the orientation of the inner ellipse depends on the laser polarization. With vertically polarized laser pulse, this ellipse orientates vertically (Fig. 2(j)). On the other hand, we observed no inner elliptic with circular polarization (Fig. 2(k)). Third, the ring-shaped radiation consists of numerous colorful needle-like rays, which are radially oriented (Fig. 1(c)-(d), Fig. 2(h)-(i)). This structure is found with linearly, circularly, and elliptically polarized laser pulses. The spectrum of this needle-like radiation is presented in Fig. 3(b), together with those of the central white-light and the incident laser pulse. The needle-like rays exhibit a similar spectrum to the central white-light, being significantly broadened compared to the incident laser.

More importantly, we found the CWL in a large number of transparent materials including 4 optical glasses (fused silica, Corning glass, BK7 glass, Pyrex), 6 optical crystals (Quartz,  $\text{CaF}_2$ ,  $\text{MgF}_2$ ,  $\text{BaF}_2$ , Sapphire, KDP) and 4 transparent polymers (polymethyl methacrylate (PMMA), polyimide (PI), polycarbonate (PC), allyl diglycol carbonate (CR-39)). For example, the result for  $\text{CaF}_2$  after 580 shots is presented in Fig. 2(l). The thickness of the various samples varies from 1.5 mm (Corning glass) to 9 mm (KDP) and we found that the CWL phenomena are similar with fused silica samples of different thickness (3 mm and 7 mm). However, the appearance of the ellipse radiation is sensitive to the nature of materials and the laser energy, a fact that requires a future systematic study. We therefore mainly focus on the CWL in this article. We also demonstrated the generality of CWL with respect to laser pulse duration and focusing geometry. For pulses of 500 fs and 1 ps duration with sufficient energy, we observed similar phenomenon as presented in Fig. 2. We changed the 1000 mm lens to 500, 200, 100, 50, 25 mm and always observed CWL emission from fused silica sample.

In the following, we propose an explanation for this spectacular phenomenon. We noticed that visible surface damage on the sample is produced once the incident laser energy exceeds the threshold for CWL generation. This indicates the CWL is closely related to the surface damage. For examination of the damage morphology, the sample is observed with a scanning electron microscope (JEOL JSM 6300) after being cleaned in ultrasonic methanol bath and coated for better surface conductivity. Figure 4(a) presents the damage morphology corresponding to  $N = 240$  shots for horizontal laser polarization. There are two regimes with distinct features. Scanning electron microscope (SEM) images of those regimes are shown in Fig. 4(b) and Fig. 4(c). Well-aligned periodic ripple structures cover the central area (regime I), while the periphery is covered by irregularly distributed submicron rods and particles (regime II). Laser induced periodic surface structures (LIPSS) have attracted growing research interests since their discovery 40 years ago [16-19]. Here the ripple structure is parallel to the laser polarization (Fig. 4(b)), referred to as c-type ripple in the literature [16, 17].

We systematically examined laser damage on fused silica as a function of pulse numbers. It was found that after 5 shots, the damage site is entirely covered by the LIPSS. With increasing laser shots, the LIPSS on the peripheral area disappears progressively towards the center while it transforms into regime II. Further laser shots finally lead to disappearance of the central LIPSS. Figure 4(d) shows the dependence of the central and peripheral damage areas on the number of laser shots.

The simultaneous disappearance of the central LIPSS and of the inner ellipse-shaped pattern around 400 shots (Fig. 2(g)-(h) and Fig. 4(d)) suggests that the latter originates from the LIPSS. In fact, a similar effect was observed in a reflected geometry on metal surfaces with a continuous probe laser [16]. This naturally explains the polarization dependence of the orientation of the inner elliptic radiation. The underlying physical mechanism is that the direction of the LIPSS depends on the laser polarization [16-19]. For circularly polarized laser, no well defined LIPSS is generated. This agrees with the fact that no inner ellipse-shaped radiation is recorded in the corresponding case (Fig. 2(k)).

We further examined the damage profiles with a profilometer (Dektak 150). The results are presented in Fig. 4(e). The damage craters deepen progressively with successive laser shots. After about 400 shots the depth of crater becomes saturated, mainly due to the reduction of laser intensity caused by the increase of the lateral surface of the crater irradiated by the laser pulse. In the presence of the crater structure, the laser beam is expected to diffract. Considering the model of a conical crater sketched in Fig. 5(a), the relationship between the opening angle  $2\phi_s$  of a collimated incident beam and the crater apex angle  $\phi_1$  can be obtained as:

$$n_0 \sin(\phi_s) = n_1 \sin\left(\frac{\pi - \phi_1}{2} - \arcsin\left(\frac{n_0 \sin((\pi - \phi_1)/2)}{n_1}\right)\right). \quad (1)$$

Here  $n_0$  and  $n_1$  are the refractive index of air and the fused silica sample respectively. For quantitative analysis, we plot in Fig. 4(f) the measured apex angle  $\phi_1$  of the crater, as well as the corresponding calculated opening angle  $2\phi_s$  based on Eq. (1). The calculated results agree well with the measured opening angle of the CWL (Fig. 3(a)), indicating that the opening of the CWL is due to the gradual shape transformation of the crater under successive laser shots.

Based on these understandings, we performed numerical simulations of the beam propagation through the sample to obtain insight into the formation of such needle-like structure of the CWL. We take a spatial Gaussian beam  $A(x, y)$  containing a flat spectrum ranging from 500 nm to 1000 nm as incident laser, for simplicity. The measured crater spatial profile  $P(x, y)$ , and the SEM surface morphology  $S(x, y)$  corresponding to the stable state ( $N > 500$ ) are considered to impart spatial phases on the incident beam. For the sake of numerical feasibility, we obtained the radiation pattern at infinity by taking Fourier transformation of each incident spectrum component modulated by these spatial phases. By adding up the radiation patterns of each spectrum component assigned a false-color, the total radiation pattern in the far field  $F(k_x, k_y)$  is obtained as

$$F(k_x, k_y) = \sum_{\lambda=500nm}^{1000nm} \iint A(x, y) \exp\{-ik(n_1 - n_0)(P(x, y) + S(x, y) \cdot \Delta d)\} \exp[-i(k_x x + k_y y)] dx dy, \quad (2)$$

where  $k = 2\pi/\lambda$  is the wavenumber in air and  $\Delta d$  is the modulation depth of the ablated surface. The opening angle is related to the wavevector as  $\theta_{x,y} = \arcsin(n_1 k_{x,y}/k)$ . In the simulation, we tested different modulation depths and found that  $\Delta d = 1\mu m$  produced result most close to our experimental observation, which is presented in Fig. 5(b). It clearly shows the radially oriented needle-like structure, in good agreement with the measurement (Fig. 2(i)). Therefore, it confirms the CWL originates in the diffraction of the laser generated white-light on the damage crater covered by submicron structures.

Based on the above simulation, we can now explain the formation of this radially elongated needle-like feature in an intuitive manner. Due to the strong nonlinear interaction of the laser pulses and the surface, the incident femtosecond pulse evolves into a broadband white-light, as presented in Fig. 3(b). At the same time, this white-light diffracts on the damage crater covered by numerous submicron structures. As a result of diffraction, each spectrum component exhibits an identical speckle radiation pattern in the far field, except for a spatial scaling factor related to the wavelength. The superposition of a large number of identical patterns with continuously varying spatial scale naturally leads to a radially elongated needle-like structure. This is similar to the so-called ciliary corona phenomenon presented in our human eye vision system. In the night time, most people see a radiating pattern of numerous fine, slightly colored needles of light around bright lights against dark background [20]. The effect was attributed to the presence of specific distributions of small particles in natural eyes, a counter-part of the sub-micron random features on the damage crater in our experiments.

With the underlying mechanisms identified, this optical phenomenon then provides a real time *in-situ* observation of the laser ablation process. First, the appearance or not of the CWL indicates whether observable damage (by SEM)

occurs on the sample surface after successive laser irradiation. Second, the apex angle of the damage crater can be easily deduced from the opening angle of the CWL, based on Eq. (1). Finally, the species of the sub-wavelength structure (regime I and II) can be identified *in-situ* from the distribution of the white-light because the correspondence between the optical radiation pattern (CWL and the elliptic ring structure) and the nature of the damage morphology has been established.

In conclusion, we observed a universal novel phenomenon at the frontiers of nonlinear optics and material science, named as ciliary white-light, during laser ablation on transparent materials. This CWL features as a ring-shaped irradiation consist of numerous needle-like rays which are radially orientated. In addition to the CWL, another polarization-dependent ellipse-shaped radiation pattern was observed. We found that the CWL emission is generic with respect to laser parameters and the nature of materials while the ellipse-shaped radiation is not. We interpret the formation of CWL and the pulse evolution in terms of nonlinear diffraction of the laser pulse on the damage crater covered with different kinds of laser-induced surface nanostructures. Calculations and simulations agree well with the observations. This optical phenomenon carries rich structural information of the surface damage, providing a real time *in-situ* observation of the laser ablation process.

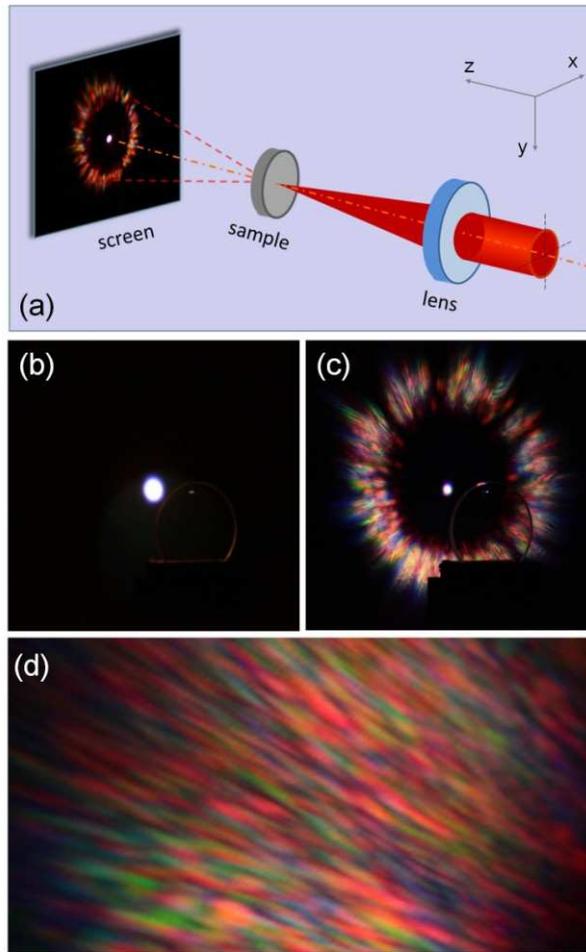
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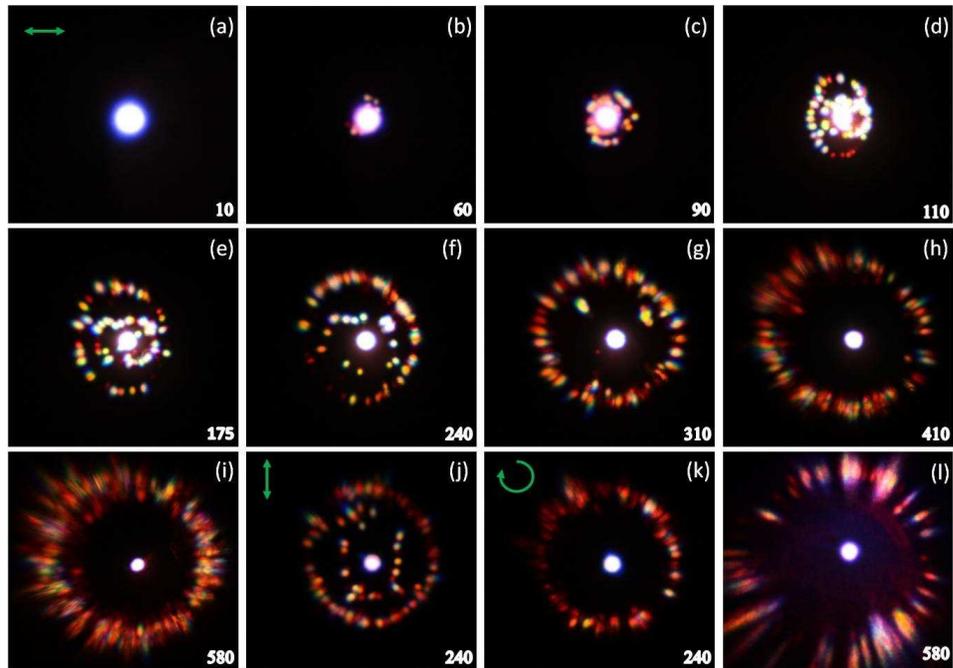
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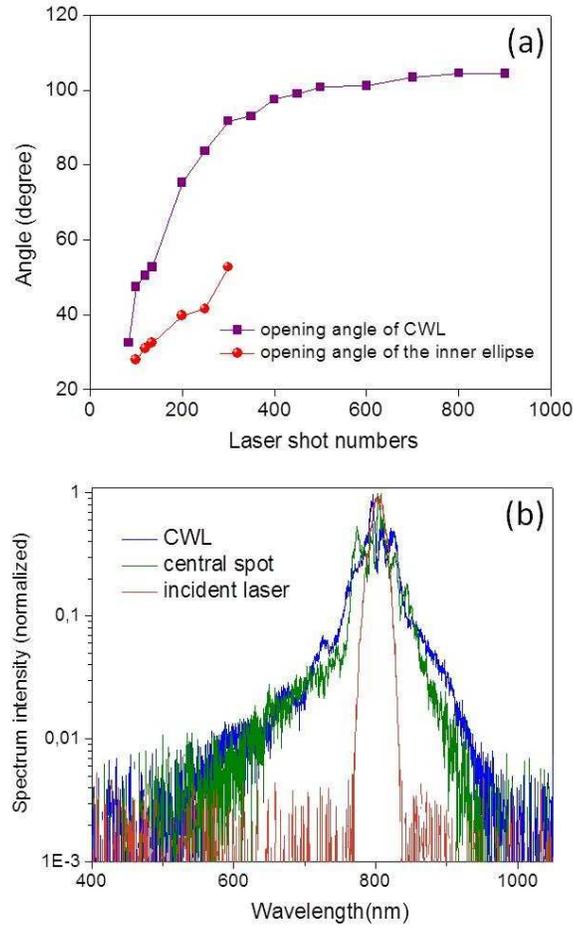
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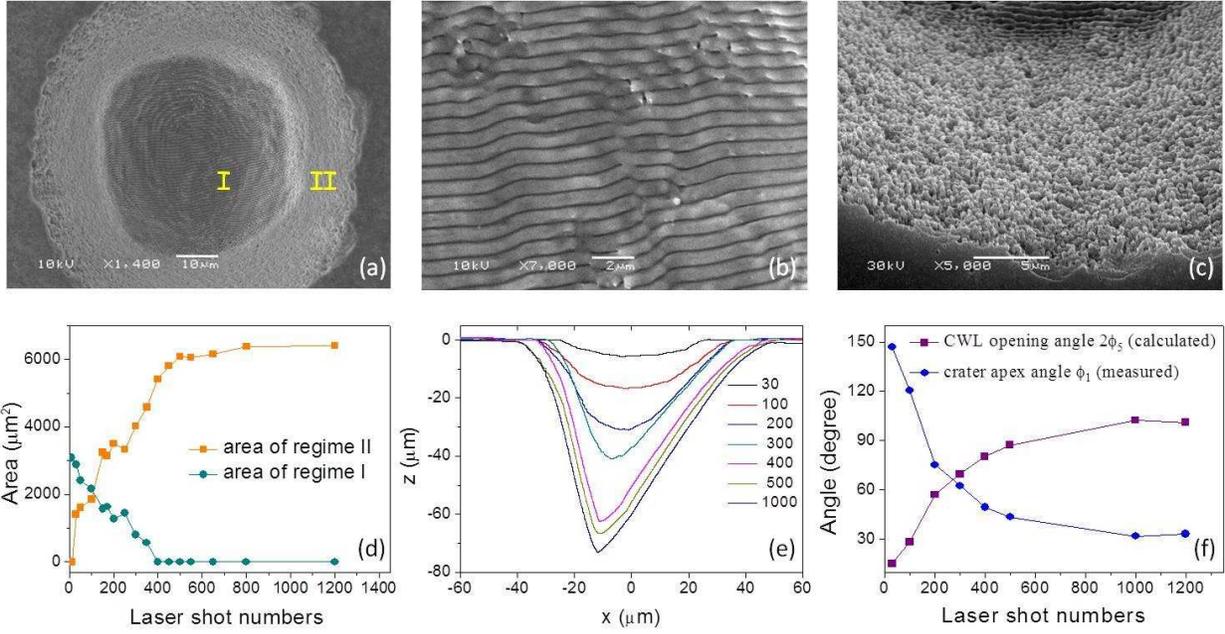
**Fig. 1.** (a), Experiment setup. (b), White-light generated by pulses of energy  $E_m = 15 \mu\text{J}$ . (c), Ciliary white-light generated by pulses of  $210 \mu\text{J}$ . In (b) and (c), scattering on the front surface of the sample is also visible. (d), Magnified structure of the CWL at the top-left section.



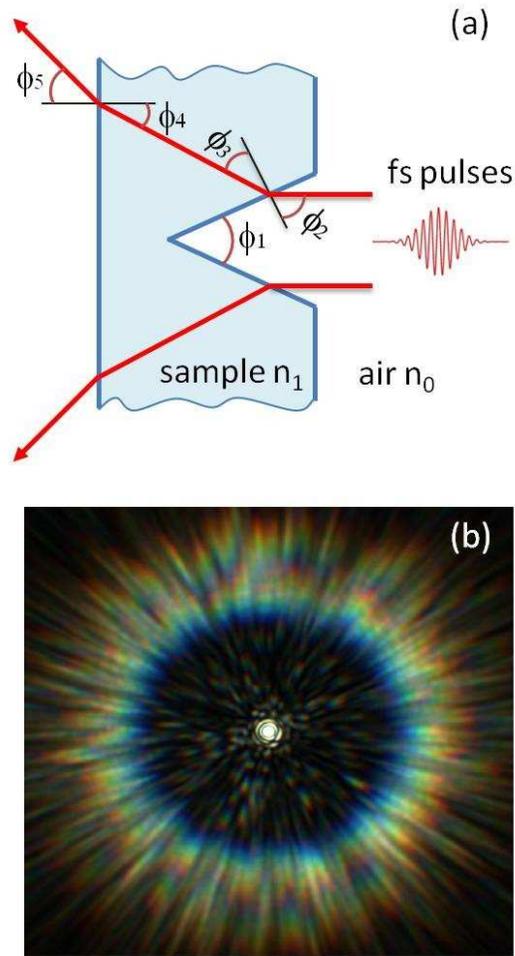
**Fig. 2.** (a)-(i), Laser emission patterns as a function of laser shot numbers for horizontally polarized pulses. The number of laser shots is indicated in each panel. The laser pulse ( $210 \mu\text{J}$ ) is focused by a lens with  $f = 1000 \text{ mm}$ . (j) and (k), Laser emission patterns for 240 laser shots with vertical and circular laser polarizations. Other conditions are the same as in (a). (l), Emission pattern obtained with  $\text{CaF}_2$  for 580 laser shots, with other experimental conditions identical to (a). The full view angle of the image is  $120^\circ \times 117.5^\circ$  with respect to the laser ablation spot on the sample, for horizontal and vertical directions.



**Fig. 3.** (a), Opening angles of the CWL and the inner ellipse-shaped radiation (vertical direction) as a function of laser shot numbers. (b), Spectra of the CWL, the central white-light, and the incident laser pulses. The CWL spectrum is an average over 10 different positions on the ring structure.



**Fig. 4.** (a), Morphology of the damage produced with 240 laser shots observed with SEM. Regimes I and II represent the central area covered by LIPSS and the peripheral area covered by nanostructures respectively. (b) and (c), Higher resolution SEM image of the centre of regime I and the bottom of the regime II. (d), Evolution of the area of the two regimes. (e), Profiles of the damage produced with successive laser shots. (f), Apex angle  $\phi$  determined from (e) and the corresponding calculated opening angle  $2\phi_s$  of a collimated beam shot at the damage crater.



**Fig. 5.** (a), Modeling of the damage crater as a cone and its consequence on an incident collimated beam. (b), Simulated result, to be compared with the experimental result in Fig. 2(i). The view angle is  $120^\circ \times 120^\circ$ .