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ON THE CLASSICAL LIMIT OF THE SCHRÖDINGER EQUATION

CLAUDE BARDOS, FRANÇOIS GOLSE, PETER MARKOWICH, AND THIERRY PAUL

ABSTRACT. This paper provides an elementary proof of the classical limit of the Schrödinger equation with WKB type initial data and over arbitrary long finite time intervals. We use only the stationary phase method and the Laptev-Sigal simple and elegant construction of a parametrix for Schrödinger type equations [A. Laptev, I. Sigal, *Review of Math. Phys.* **12** (2000), 749–766]. We also explain in detail how the phase shifts across caustics obtained when using the Laptev-Sigal parametrix are related to the Maslov index.

1. THE CLASSICAL SCALING

Consider the evolution Schrödinger equation

$$i\hbar\partial_t\psi = -\frac{\hbar^2}{2m}\Delta_x\psi + V(x)\psi$$

for the wave function ψ of a point particle of mass m subject to the action of an external potential $V \equiv V(x) \in \mathbf{R}$.

Choosing “appropriate” units of time T and length L , we recast the Schrödinger equation in terms of dimensionless variables $\hat{t} := t/T$ and $\hat{x} := x/L$. We define a rescaled wave function $\hat{\psi}$ and a rescaled, dimensionless potential \hat{V} by the formulas

$$\hat{\psi}(\hat{t}, \hat{x}) := \psi(t, x) \quad \text{and} \quad \hat{V}(\hat{x}) := \frac{T^2}{mL^2}V(x).$$

In these dimensionless variables, the Schrödinger equation takes the form

$$i\partial_{\hat{t}}\hat{\psi} = -\frac{\hbar T}{2mL^2}\Delta_{\hat{x}}\hat{\psi} + \frac{mL^2}{\hbar T}\hat{V}(\hat{x})\hat{\psi}.$$

The dimensionless number $2\pi\hbar T/mL^2$ is the ratio of the Planck constant to mL^2/T , that is (twice) the action of a classical particle of mass m moving at speed L/T on a distance L . If the scales of time T and length L have been chosen conveniently, L/T is the typical order of magnitude of the particle speed, and L is the typical length scale on which the particle motion is observed. The classical limit of quantum mechanics is defined by the scaling assumption $2\pi\hbar \ll mL^2/T$ — i.e. the typical action of the particle considered is large compared to \hbar . Equivalently, mL/T is the order of magnitude of the particle momentum, so that $2\pi\hbar T/mL$ is its de Broglie wavelength; the scaling assumption $2\pi\hbar T/mL \ll L$ means that the de Broglie wavelength of the particle under consideration is small compared to the observation length scale L .

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Introducing the small, dimensionless parameter $\epsilon = \hbar T/mL^2$ and dropping hats in the dimensionless variables as well as on the rescaled wave function and dimensionless potential, we arrive at the following formulation for the Schrödinger equation in dimensionless variables

$$(1) \quad i\epsilon \partial_t \psi = -\frac{1}{2}\epsilon^2 \Delta_x \psi + V(x)\psi.$$

The WKB ansatz postulates that, at time $t = 0$, the wave function is of the form

$$\psi(t, x) = a^{in}(x) e^{iS^{in}(x)/\hbar}, \quad x \in \mathbf{R}^N.$$

Consistently with the scaling argument above, we set

$$\hat{a}^{in}(\hat{x}) := a^{in}(x) \text{ and } \hat{S}^{in}(\hat{x}) := TS^{in}(x)/mL^2$$

— since S^{in} has the dimension of an action — so that

$$(2) \quad \hat{\psi}(0, \hat{x}) = \hat{a}^{in}(\hat{x}) e^{i\hat{S}^{in}(\hat{x})/\epsilon}.$$

Dropping hats in the initial data as well as in the Schrödinger equation, one arrives at the following Cauchy problem for the Schrödinger equation in dimensionless variables:

$$(3) \quad \begin{cases} i\epsilon \partial_t \psi_\epsilon = -\frac{1}{2}\epsilon^2 \Delta_x \psi_\epsilon + V(x)\psi_\epsilon, & x \in \mathbf{R}^N, t \in \mathbf{R}, \\ \psi_\epsilon(0, x) = a^{in}(x) e^{iS^{in}(x)/\epsilon}. \end{cases}$$

The problem of the classical limit of the Schrödinger equation is to describe the wave function ψ_ϵ in the limit as $\epsilon \rightarrow 0^+$.

The main result obtained in this paper (Theorem 4.1) is of course not new — it is stated without proof as Theorem 5.1 in [2]; see also Example 6.1 described on pp. 141-143 in [15], especially formula (1.21) there.

Yet our purpose in the present work is to provide a short and self-contained proof of this result, based on Laptev-Sigal lucid construction of a parametrix for the Schrödinger equation (1) in [12]. In this way, we avoid using either Maslov's formalism [15, 16] of the canonical operator, or the semiclassical analogue of Hörmander's global theory [11] of Lagrangian distributions or Fourier integral operators. The Laptev-Sigal construction, which is remarkably short (pp. 753–759 in [12]), uses instead Fourier integrals with a complex phase whose imaginary part is quadratic.

2. THE CLASSICAL DYNAMICS

Before discussing the classical limit of the Schrödinger equation, we need to recall a few preliminary results on either the quantum dynamics defined by (3) or the classical dynamics in phase space associated to the Schrödinger equation (3).

Assume that $V \in C^\infty(\mathbf{R}^N)$ satisfies

$$(4) \quad \partial^\alpha V \in L^\infty(\mathbf{R}^N) \quad \text{for each multi-index } \alpha \in \mathbf{N}^N \text{ such that } |\alpha| > 0$$

and

$$(5) \quad \frac{V(x)}{|x|} \rightarrow 0 \quad \text{as } |x| \rightarrow +\infty.$$

Then the Hamiltonian

$$H(x, \xi) := \frac{1}{2}|\xi|^2 + V(x)$$

generates a global flow

$$\mathbf{R}^N \times \mathbf{R}^N \ni (x, \xi) \mapsto \Phi_t(x, \xi) := (X_t(x, \xi), \Xi_t(x, \xi)) \in \mathbf{R}^N \times \mathbf{R}^N$$

that is of class C^1 . This Hamiltonian flow is the classical dynamics corresponding to the quantum dynamics defined by the Schrödinger equation (3).

This classical dynamics satisfies the following properties: for each $\eta > 0$, there exists $C_\eta > 0$ such that

$$\sup_{|t| \leq T} |X_t(x, \xi) - x| \leq C_\eta(1 + |\xi|) + \eta|x|$$

for each $x, \xi \in \mathbf{R}^N$, and

$$|D\Phi_t(x, \xi) - \text{Id}_{\mathbf{R}^N \times \mathbf{R}^N}| \leq e^{\kappa|t|} - 1$$

for all $t \in \mathbf{R}$. (See Lemma 4.1 in [5].)

Assume further that

$$(6) \quad \sup_{x \in \mathbf{R}^N} \int_{\mathbf{R}^N} \Gamma_\eta(x-y)V^-(y)dy \rightarrow 0 \text{ as } \eta \rightarrow 0 \quad \text{if } N \geq 2$$

with

$$\Gamma_\eta(z) = \begin{cases} \mathbf{1}_{[0, \eta]}(|z|)|z|^{2-N} & \text{if } N \geq 3, \\ \mathbf{1}_{[0, \eta]}(|z|)\ln(1/|z|) & \text{if } N = 2, \end{cases}$$

while

$$(7) \quad \sup_{x \in \mathbf{R}^N} \int_{x-1}^{x+1} V^-(y)dy < \infty \quad \text{if } N = 1.$$

Under assumptions (6)-(7), the operator $-\frac{1}{2}\epsilon^2\Delta_x + V$ has a self-adjoint extension on $L^2(\mathbf{R}^N)$ that is bounded from below. (See [14] on p. 567.)

Next we discuss the properties of the initial phase function S^{in} . It will be convenient to assume that S^{in} is of class C^2 at least on \mathbf{R}^N and satisfies the following growth condition at infinity:

$$(8) \quad \frac{|\nabla S^{in}(y)|}{|y|} \rightarrow 0 \quad \text{as } |y| \rightarrow 0.$$

Consider the map

$$(9) \quad F_t : \mathbf{R}^N \ni y \mapsto F_t(y) = X_t(y, \nabla S^{in}(y)) \in \mathbf{R}^N,$$

together with the absolute value of its Jacobian determinant

$$(10) \quad J_t(y) = |\det(DF_t(y))|.$$

We also introduce the set

$$(11) \quad \begin{aligned} C &:= \{(t, x) \in \mathbf{R} \times \mathbf{R}^N \mid F_t^{-1}(\{x\}) \cap J_t^{-1}(\{0\}) \neq \emptyset\}, \\ C_t &:= \{x \in \mathbf{R}^N \mid (t, x) \in C\}, \quad t \in \mathbf{R}. \end{aligned}$$

For lack of a better terminology and by analogy with geometric optics, C will be referred to as the “caustic” set, and C_t as the “caustic fiber”.

Proposition 2.1. *Assume that (8) holds, and that*

$$|\nabla^2 S^{in}(y)| = O(1) \quad \text{as } |y| \rightarrow +\infty.$$

(a) *For each $t \in \mathbf{R}$, the map F_t is proper and onto. More precisely, for each $T > 0$*

$$(12) \quad \sup_{|t| \leq T} \frac{|F_t(y) - y|}{|y|} \rightarrow 0 \quad \text{as } |y| \rightarrow \infty.$$

(b) *The caustic set C is closed in $\mathbf{R} \times \mathbf{R}^N$, and $\mathcal{L}^N(C_t) = 0$ for each $t \in \mathbf{R}$.*

(c) For each $(t, x) \in \mathbf{R} \times \mathbf{R}^N \setminus C$, the set $F_t^{-1}(\{x\})$ is finite, henceforth denoted by

$$\{y_j(t, x), j = 1, \dots, \mathcal{N}(t, x)\}.$$

The integer \mathcal{N} is a constant function of (t, x) in each connected component of $\mathbf{R} \times \mathbf{R}^N \setminus C$ and, for each $j \geq 1$, the map y_j is of class C^1 on each connected component of $\mathbf{R} \times \mathbf{R}^N \setminus C$ where $\mathcal{N} \geq j$.

(d) There exists $a < 0 < b$ such that $C \cap ((a, b) \times \mathbf{R}^N) = \emptyset$ and $\mathcal{N} = 1$ on $(a, b) \times \mathbf{R}^N$.

(e) For each $(t, x) \in \mathbf{R} \times \mathbf{R}^N \setminus C$, the integer $\mathcal{N}(t, x)$ is odd.

(See Theorems 2.3 and 2.5 in [5] for a proof of these results.)

One might be worried that the sublinearity assumption (5) obviously excludes the harmonic oscillator, i.e. the case where $V(x) := \frac{1}{2}|x|^2$. In view of the importance of the harmonic oscillator in quantum mechanics, the setting presented in this section might seem unfelicitous. In fact, the main reason for assuming (5) is that the set $F_t(\{x\})$, which consists of isolated points only, is necessarily finite because F_t is proper so that $F_t(\{x\})$ is a compact set with only isolated points, i.e. a finite set. For that reason, our main results Theorems 3.1 and 4.1 hold without additional assumption on the initial amplitude a^{in} . If (5) is not satisfied, one cannot in general conclude that $F_t(\{x\})$ is finite. But one can choose to consider only the case of initial amplitudes a^{in} with compact support, so that $F_t(\{x\}) \cap \text{supp}(a^{in})$ is necessarily finite. With this slight modification, most of the results of the present paper also apply to the harmonic oscillator.

3. THE FREE CASE $V \equiv 0$.

For the sake of clarity, we first discuss in detail the case $V \equiv 0$. This section is the exact analogue for the Schrödinger equation of chapter XII.1 in [9], which presents geometric optics as an asymptotic theory for solutions of the wave equation with constant coefficients in the high frequency limit.

Assume that $a^{in} \in C_c(\mathbf{R}^N)$ while $S^{in} \in C^2(\mathbf{R}^N)$ and satisfies (8). The solution ψ_ϵ of (3) is given by the explicit formula

$$(13) \quad \psi_\epsilon(t, x) = \left(\frac{1}{\sqrt{2\pi i \epsilon t}} \right)^N \int_{\mathbf{R}^N} e^{i\phi(t, x, y)/\epsilon} a^{in}(y) dy, \quad x \in \mathbf{R}^N, t \neq 0,$$

where $\sqrt{\cdot}$ designates the holomorphic extension of the square root to $\mathbf{C} \setminus \mathbf{R}_-$ and

$$(14) \quad \phi(t, x, y) := \frac{|x - y|^2}{2t} + S^{in}(y), \quad x, y \in \mathbf{R}^N, t \neq 0.$$

Applying the stationary phase method (see Theorems 7.7.5 and 7.7.6 in [8]) to the explicit formula (13) leads to the following classical result on the asymptotic behavior of ψ_ϵ as $\epsilon \rightarrow 0^+$.

Theorem 3.1. *Let $S^{in} \in C^{m+1}(\mathbf{R}^N)$ satisfy (8) and $a^{in} \in C_c^m(\mathbf{R}^N)$ with $m \geq \frac{3}{2}N + 3$. For all $(t, x) \in \mathbf{R} \times \mathbf{R}^N \setminus C$ and $\epsilon > 0$, the solution ψ_ϵ of the Cauchy problem (3) satisfies*

$$(15) \quad \psi_\epsilon(t, x) - \sum_{j=1}^{\mathcal{N}(t, x)} \frac{i^{-\mathcal{M}_j(t, x)}}{|\det(I + t\nabla^2 S^{in}(y_j(t, x)))|^{1/2}} a^{in}(y_j(t, x)) e^{i\phi(t, x, y_j(t, x))/\epsilon} = O(\epsilon)$$

uniformly as (t, x) runs through compact subsets of $\mathbf{R} \times \mathbf{R}^N \setminus C$, where $\mathcal{M}_j(t, x)$ is the number of negative eigenvalues of the matrix $I + t\nabla^2 S^{in}(y_j(t, x))$.

Proof. Pick $(t, x) \in \text{supp}(a^{in})$; critical points y of the phase ϕ are obtained by solving for y the equation

$$\nabla_y \phi(t, x, y) = \frac{y - x}{t} + \nabla S^{in}(y) = 0,$$

i.e.

$$x = y + t\nabla S^{in}(y) := F_t(y).$$

If $(t, x) \notin C$, the equation above has finitely many solutions denoted

$$y_1(t, x), \dots, y_{\mathcal{N}(t, x)}$$

(see Proposition 2.1 above). For $j = 0, \dots, \mathcal{N}(t, x)$, let $\chi_j \in C_c^\infty(\mathbf{R}^N)$ satisfy

$$\chi_j \equiv 1 \text{ near } y_j(t, x), \quad j = 1, \dots, \mathcal{N}(t, x), \quad F_t^{-1}(\{x\}) \cap \text{supp}(\chi_0) = \emptyset$$

and

$$\chi_j \geq 0 \text{ for all } j = 0, \dots, \mathcal{N}(t, x), \quad \sum_{j=0}^{\mathcal{N}(t, x)} \chi_j \equiv 1 \text{ near } \text{supp}(a^{in}).$$

Thus

$$(16) \quad \psi_\epsilon(t, x) = \left(\sqrt{2\pi i \epsilon t}\right)^{-N} \sum_{0 \leq j \leq \mathcal{N}(t, x)} \int_{\mathbf{R}^N} e^{i\phi(t, x, y)/\epsilon} a^{in}(y) \chi_j(y) dy.$$

In the term corresponding with $j = 0$, the phase $\phi(t, x, \cdot)$ has no critical point on $\text{supp}(\chi_0 a^{in})$. By Theorem 7.7.1 in [8], one has

$$(17) \quad \left| \int_{\mathbf{R}^N} e^{i\phi(t, x, y)/\epsilon} a^{in}(y) \chi_0(y) dy \right| \leq (\epsilon|t|)^k C_0(t) \|\nabla^{k+1} S^{in}\|_{L^\infty(B(0, R))} \sum_{|\alpha| \leq k} \sup_{y \in \mathbf{R}^N} \frac{|\partial^\alpha (a^{in} \chi_0)(y)|}{|y + t\nabla S^{in}(y) - x|^{|\alpha| - 2k}},$$

where $R > 0$ is chosen so that $\text{supp}(a^{in}) \subset B(0, R)$ and the function C_0 is bounded on bounded sets of \mathbf{R}_+ . Thus

$$\left| \left(\sqrt{2\pi i \epsilon t}\right)^{-N} \int_{\mathbf{R}^N} e^{i\phi(t, x, y)/\epsilon} a^{in}(y) \chi_0(y) dy \right| = O(\epsilon|t|)$$

provided that one can choose $k \geq \frac{1}{2}N + 1$ in the previous inequality. Applying Theorem 7.7.1 of [8] for this estimate requires $m \geq k \geq \frac{1}{2}N + 1$.

For $j = 1, \dots, \mathcal{N}(t, x)$, one has

$$\begin{aligned} & \left| \int_{\mathbf{R}^N} e^{i\phi(t,x,y)/\epsilon} a^{in}(y) \chi_j(y) dy - (2\pi\epsilon)^{N/2} \frac{(a^{in} \chi_j)(y_j(t, x)) e^{i\phi(t,x,y_j(t,x))/\epsilon}}{\sqrt{(-i)^N \det \nabla_y^2 \phi(t, x, y_j(t, x))}} \right| \\ & \leq (\epsilon|t|)^k C_1 \left(\sup_{\chi_j(y) > 0} \frac{t|y - y_j(t, x)|}{|y + t\nabla S^{in}(y) - x|} + t \|\nabla^{3k+1} S^{in}\|_{L^\infty(B(0,R))} \right) \\ & \quad \times \sum_{|\alpha| \leq k} \sup_{y \in \mathbf{R}^N} |\partial^\alpha (a^{in} \chi_0)(y)| \\ & \quad + (2\pi\epsilon)^{N/2} \frac{|(a^{in} \chi_j)(y_j(t, x))|}{|\det \nabla_y^2 \phi(t, x, y_j(t, x))|^{1/2}} \sum_{1 \leq j < k} (\epsilon|t|)^j L_j a^{in} \end{aligned}$$

by Theorem 7.7.5 in [8], where the function C_1 is bounded on bounded sets of \mathbf{R}_+ . In the inequality above, we have denoted

$$L_j a^{in} := \sum_{\substack{\nu - \mu = j \\ 2\nu \geq 3\mu}} \frac{1}{2^\nu \mu! \nu!} |((I + t\nabla^2 S^{in}(y_j(t, x)))^{-1} : \nabla_y^{\otimes 2})^\nu (a^{in} T_{y_j}^2 \phi(t, x, \cdot))(y_j(t, x))|,$$

and

$$T_y^2 f(x) := f(x) - f(y) - f'(y) \cdot (x - y) - \frac{1}{2} f''(y) \cdot (x - y)^2.$$

Thus, Theorem 7.7.5 of [8] shows that

$$(18) \quad \left| \left(\sqrt{2\pi i \epsilon t} \right)^{-N} \int_{\mathbf{R}^N} e^{i\phi(t,x,y)/\epsilon} a^{in}(y) \chi_j(y) dy - \frac{a^{in}(y_j(t, x)) e^{i\phi(t,x,y_j(t,x))/\epsilon - iN\pi/4}}{\sqrt{(-i)^N \det(I + t\nabla^2 S^{in}(y_j(t, x))}} \right| = O(\epsilon|t|)$$

provided that one can choose $k \geq \frac{1}{2}N + 1$ in the previous inequality. This requires $m \geq 3k \geq \frac{3}{2}N + 3$.

It remains to identify $\mathcal{M}_j(t, x)$. The matrix $\nabla^2 S^{in}(y)$ is symmetric with real entries, and therefore reducible to diagonal form with real eigenvalues

$$\lambda_1(y) \geq \lambda_2(y) \geq \dots \geq \lambda_N(y)$$

counted with multiplicities. Since $(t, x) \notin C$, one has $1 + t\lambda_k(y_j(t, x)) \neq 0$ for each $k = 1, \dots, N$ and each $j = 1, \dots, \mathcal{N}(t, x)$. Set

$$\mathcal{M}_j(t, x) := \#\{1 \leq k \leq N \text{ s.t. } 1 + t\lambda_k(y_j(t, x)) < 0\}.$$

Thus

$$\det(I + t\nabla^2 S^{in}(y_j(t, x))) = (-1)^{\mathcal{M}_j(t, x)} |\det(I + t\nabla^2 S^{in}(y_j(t, x)))|,$$

and therefore

$$\begin{aligned} & \sqrt{(-i)^N \det(I + t\nabla^2 S^{in}(y_j(t, x)))} \\ & = |\det(I + t\nabla^2 S^{in}(y_j(t, x)))|^{1/2} \sqrt{-i}^{N - \mathcal{M}_j(t, x)} \sqrt{i}^{\mathcal{M}_j(t, x)} \\ & = |\det(I + t\nabla^2 S^{in}(y_j(t, x)))|^{1/2} e^{-i\pi(N - \mathcal{M}_j(t, x) - \mathcal{M}_j(t, x))/4}. \end{aligned}$$

Thus

$$\begin{aligned}
 (19) \quad & \frac{a^{in}(y_j(t, x))e^{i\phi(t, x, y_j(t, x))/\epsilon - iN\pi/4}}{\sqrt{(-i)^N \det(I + t\nabla^2 S^{in}(y_j(t, x)))}} \\
 &= \frac{a^{in}(y_j(t, x))e^{i\phi(t, x, y_j(t, x))/\epsilon - iN\pi/4} e^{+i\pi(N-2\mathcal{M}_j(t, x))/4}}{|\det(I + t\nabla^2 S^{in}(y_j(t, x)))|^{1/2}} \\
 &= \frac{a^{in}(y_j(t, x))e^{i\phi(t, x, y_j(t, x))/\epsilon - \mathcal{M}_j(t, x)}}{|\det(I + t\nabla^2 S^{in}(y_j(t, x)))|^{1/2}}.
 \end{aligned}$$

The proof of (15) follows from substituting (19) in (18), and using (17) and (18) in the sum (16) representing ψ_ϵ . \square

Several remarks are in order.

First, even if S^{in} does not satisfy the assumption (8), for $(t, x) \in \mathbf{R} \times \mathbf{R}^N \setminus C$, the set $F_t^{-1}(\{x\})$ of critical points of the map $y \mapsto \phi(t, x, y)$ is discrete by the implicit function theorem. Its intersection with $\text{supp}(a^{in})$ is therefore finite whenever a^{in} has compact support.

Another observation bears on the physical meaning of the points $y_j(t, x)$. These points satisfy

$$y_j(t, x) + t\nabla S^{in}(y_j(t, x)) = x.$$

In other words, x is the position at time t of the free particle leaving the position $y_j(t, x)$ at time $t = 0$ with velocity $\nabla S^{in}(y_j(t, x))$, according to Newton's first law of classical mechanics.

Next we recall the following well known geometric interpretation of the caustic C . In the case $N = 1$, the system (11) reduces to

$$\begin{cases} y + t(S^{in})'(y) = x, \\ 1 + t(S^{in})''(y) = 0. \end{cases}$$

The first equality is the defining equation for a family of straight lines in $\mathbf{R}_t \times \mathbf{R}_x$ parametrized by $y \in \mathbf{R}$ — the trajectories of the particle in classical mechanics — while the second equality follows from deriving the first with respect to the parameter y . By eliminating the parameter y between both equations, we see that the caustic C is in general the envelope of the family of straight lines defined by the first equation.

The integer $\mathcal{M}_j(t, x)$ can be interpreted in terms of the notion of Maslov index, as explained in the last section of this paper. Notice that the integer $\mathcal{M}_j(t, x) \rightarrow 0$ as $t \rightarrow 0$, so that, for each $x \in \mathbf{R}^N$, one has $\mathcal{M}_j(t, x) = 0$ for all t near 0. Furthermore, \mathcal{M}_j is a locally constant function of $(t, x) \in \mathbf{R} \times \mathbf{R}^N \setminus C$, and therefore the presence of the Maslov index in (15) is equivalent to a phase shift of an integer multiple of $\pi/2$ whenever (t, x) is moved across C from one connected component of $\mathbf{R} \times \mathbf{R}^N$ to the next.

4. THE CASE $V \neq 0$

The explicit representation formula (13) obviously is a major ingredient in the proof of Theorem 3.1. In the case $V \neq 0$, there is no explicit formula analogous to (13) giving the solution of the Cauchy problem for the Schrödinger equation in general.

However, under assumption (4), there exists a FIO that is a parametrix for the operator

$$G_\epsilon(t) := e^{i\frac{t}{\epsilon}\left(\frac{1}{2}\epsilon^2\Delta_x - V\right)}.$$

Perhaps the simplest approach to this most important result is Theorem 2.1 in [12], whose main features are recalled below.

Consider the action

$$(20) \quad S(t, x, \xi) := \int_0^t \left(\frac{1}{2} |\Xi_s(x, \xi)|^2 - V(X_s(x, \xi)) \right) ds$$

Given $T > 0$, we shall have to deal with the class of phase functions

$$\varphi \equiv \varphi(t, x, y, \eta) \in C \text{ of class } C^\infty \text{ on } [0, T] \times \mathbf{R}^N \times \mathbf{R}^N \times \mathbf{R}^N$$

satisfying the conditions

$$(21) \quad \begin{cases} \varphi(t, X_t(y, \eta), y, \eta) = S(t, y, \eta), \\ D_x \varphi(t, X_t(y, \eta), y, \eta) = \Xi_t(y, \eta), \\ iD_x^2 \varphi(t, x, y, \eta) \leq 0 \text{ is independent of } x, \\ \det(D_{x\eta} \varphi(t, X_t(y, \eta), y, \eta)) \neq 0 \text{ for each } (t, y, \eta) \in [0, T] \times \mathbf{R}^N \times \mathbf{R}^N. \end{cases}$$

Pick $\chi \in C_c^\infty(\mathbf{R}^N \times \mathbf{R}^N)$ and $T > 0$. Then, for any phase function φ satisfying (21) and any $n \geq 0$, there exists $A_n \equiv A_n(t, y, \eta, \epsilon) \in C_c^\infty([0, T] \times \mathbf{R}^N \times \mathbf{R}^N)[\epsilon]$ such that the FIO $G_n(t)$ with Schwartz kernel

$$(22) \quad G_{\epsilon, n}(t, x, y) = \int A_n(t, y, \eta, \epsilon) e^{i\varphi(t, x, y, \eta)/\epsilon} \frac{d\eta}{(2\pi\epsilon)^N}$$

satisfies

$$(23) \quad \sup_{0 \leq t \leq T} \|(G_\epsilon(t) - G_{\epsilon, n}(t))\chi(x, -i\epsilon\partial_x)\|_{\mathcal{L}(L^2(\mathbf{R}^N))} \leq C[V, T, \chi]\epsilon^{n-2N}.$$

In this inequality the notation $\chi(x, -i\epsilon\partial_x)$ designates the pseudo-differential operator defined by the formula

$$\chi(x, -i\epsilon\partial_x)\phi(x) := \iint_{\mathbf{R}^N \times \mathbf{R}^N} e^{i(x-y)\cdot\eta/\epsilon} \chi(x, \eta)\phi(y) \frac{dy d\eta}{(2\pi\epsilon)^N}.$$

Taking Theorem 2.1 in [12] for granted, one arrives at the following description of the classical limit of (3). It is stated without proof in Appendix 11 of [4] or as Theorem 5.1 in [2].

Let $S^{in} \in C^2(\mathbf{R}^N)$ satisfy (8) and let C be defined as in (11); let $\mathcal{N}(t, x)$ and $y_j(t, x)$ be defined as in Proposition 2.1 for each $(t, x) \in \mathbf{R} \times \mathbf{R}^N \setminus C$. Let $J_t(y)$ be defined as in (10).

Theorem 4.1. *Let $a^{in} \in C_c^m(\mathbf{R}^N)$ and $S^{in} \in C^{m+1}(\mathbf{R}^N)$ satisfy (8), with regularity index $m > 6N + 4$. For all $\epsilon > 0$ and all $(t, x) \in \mathbf{R}_+ \times \mathbf{R}^N \setminus C$, set*

$$(24) \quad \Psi_\epsilon(t, x) = \sum_{j=1}^{\mathcal{N}(t, x)} \frac{a^{in}(y_j(t, x))}{J_t(y_j(t, x))^{1/2}} e^{iS_j(t, x)/\epsilon_j - \mathcal{M}_j(t, x)},$$

with

$$S_j(t, x) := S^{in}(y_j(t, x)) + S(t, y_j(t, x), \nabla S^{in}(y_j(t, x))), \quad j = 1, \dots, \mathcal{N}(t, x),$$

where $S(t, y, \xi)$ is given by (20) and $\mathcal{M}_j(t, x) \in \mathbf{Z}$ for all $(t, x) \in \mathbf{R}_+ \times \mathbf{R}^N \setminus C$ is constant on each connected component of $\mathbf{R} \times \mathbf{R}^N$ where $j \leq \mathcal{N}$.

Then the solution ψ_ϵ of the Cauchy problem (3) satisfies

$$(25) \quad \psi_\epsilon(t, x) = \Psi_\epsilon(t, x) + R_\epsilon^1(t, x) + R_\epsilon^2(t, x)$$

for all $T > 0$, where

$$\sup_{0 \leq t \leq T} \|R_\epsilon^1\|_{L^2(B(0, R))} = O(\epsilon) \text{ for all } R > 0$$

and

$$\sup_{(t, x) \in K} |R_\epsilon^2(t, x)| = O(\epsilon) \text{ for each compact } K \subset \mathbf{R}_+ \times \mathbf{R}^N \setminus C$$

as $\epsilon \rightarrow 0^+$.

A fairly classical computation shows that each one of the phases S_j is a solution of the Hamilton-Jacobi “eikonal” equation

$$\partial_t S_j(t, x) + \frac{1}{2} |\nabla_x S_j(t, x)|^2 + V(x) = 0$$

for (t, x) belonging to the union of all the connected components of $\mathbf{R} \times \mathbf{R}^N \setminus C$ where $j \leq \mathcal{N}(t, x)$.

The integer $\mathcal{M}_j(t, x)$ in (24) is defined precisely in the proof (see formula (34) below). It turns out to be identical to the Maslov index of the path

$$[0, t] \ni s \mapsto \Phi_s(y_j(t, x), \nabla S^{in}(y_j(t, x))),$$

as explained in the last section of this paper.

A last remark bears on the formulation of the classical limit of the Schrödinger equation in terms of the Wigner transform. Given $t > 0$, assume that a^{in} has compact support with

$$F_t(\text{supp}(a^{in}) \cap C_t) = \emptyset,$$

and that the vectors $\nabla_x S_j(t, x)$ are pairwise different for $x \in F_t(\text{supp}(a^{in}))$. Denote by $W_\epsilon[\psi_\epsilon]$ the Wigner transform of ψ_ϵ , defined by the formula

$$W_\epsilon[\psi_\epsilon](t, x, \xi) = \frac{1}{(2\pi)^N} \int_{\mathbf{R}^N} e^{-i\xi \cdot y} \psi_\epsilon(t, x + \frac{1}{2}\epsilon y) \overline{\psi_\epsilon(t, x - \frac{1}{2}\epsilon y)} dy.$$

Using Theorem 4.1 above and Proposition 1.5 in [6] shows that

$$W_\epsilon[\psi_\epsilon](x, \cdot) \rightarrow \sum_{j=1}^{\mathcal{N}(t, x)} \frac{|a^{in}(y_j(t, x))|^2}{J_t(y_j(t, x))} \delta_{\nabla_x S_j(t, x)}$$

in $\mathcal{S}'(\mathbf{R}_x^N \times \mathbf{R}_\xi^N)$ as $\epsilon \rightarrow 0$. The same result can be obtained by a completely different method, avoiding the use of Theorem 4.1, and for much less regular initial phase functions S^{in} (typically, for S^{in} having second order derivatives in the Lorentz space $L_{loc}^{N,1}$): see [5] for more details on the classical limit of the Schrödinger equation with rough phase functions.

Before giving the proof of Theorem 4.1 in detail, it is worth having in mind the differences and similarities between this proof and that of Theorem 3.1. The construction of the parametrix in [12] replaces the *exact* solution of the Schrödinger equation by an *approximate* solution, whose structure is similar to the right hand side of (13), at the expense of an error estimate controlled by using the bound in (23). We conclude by applying the stationary phase method to this approximate solution exactly as in the proof of Theorem 3.1.

Notice that the reference [12] contains a proof of Theorem 4.1 in the special case where S^{in} is a *linear* function. We shall mostly follow the argument in [12], except that the case of a general phase S^{in} , not necessarily linear, requires an additional trick (see Lemma 4.2 below).

Proof. Let $\chi(x, \xi) = \chi_1(x)\chi_2(\xi)$ with $\chi_1, \chi_2 \in C_c^\infty(\mathbf{R}^N)$, satisfying

$$\mathbf{1}_{B(0,R)}(x) \leq \chi_1(x) \leq \mathbf{1}_{B(0,R+1)} \text{ and } \mathbf{1}_{B(0,Q)}(\xi) \leq \chi_2(\xi) \leq \mathbf{1}_{B(0,Q+1)}$$

for all $x, \xi \in \mathbf{R}^N$, where $R > 0$ and Q is to be chosen later. Pick $n > 2N$; then

$$\begin{aligned} \psi_\epsilon(t, \cdot) - G_{\epsilon,n}(t)\chi(x, -i\epsilon\partial_x)\psi_\epsilon^{in} &= G_\epsilon(t)(1 - \chi(x, -i\epsilon\partial_x))\psi_\epsilon^{in} \\ &\quad + (G_\epsilon(t) - G_{\epsilon,n}(t))\chi(x, -i\epsilon\partial_x)\psi_\epsilon^{in}. \end{aligned}$$

Since $G_\epsilon(t)$ is a unitary group on $L^2(\mathbf{R}^N)$

$$\begin{aligned} \|\psi_\epsilon(t, \cdot) - G_{\epsilon,n}(t)\chi(x, -i\epsilon\partial_x)\psi_\epsilon^{in}\|_{L^2(\mathbf{R}^N)} &\leq \|(1 - \chi(x, -i\epsilon\partial_x))\psi_\epsilon^{in}\|_{L^2(\mathbf{R}^N)} \\ &\quad + \|(G_\epsilon(t) - G_{\epsilon,n}(t))\chi(x, -i\epsilon\partial_x)\psi_\epsilon^{in}\|_{L^2(\mathbf{R}^N)} \\ &\leq \|(1 - \chi(x, -i\epsilon\partial_x))\psi_\epsilon^{in}\|_{L^2(\mathbf{R}^N)} \\ &\quad + C_{T,Q}\epsilon^{n-2N}\|a^{in}\|_{L^2(\mathbf{R}^N)} \end{aligned}$$

for all $t \in [0, T]$, where $C_{Q,T} = C[V, T, \chi]$.

Now, $\chi(x, -i\epsilon\partial_x)\psi_\epsilon^{in} = \chi_1(x)\chi_2((-i\epsilon\partial_x)\psi_\epsilon^{in})$ and since $\text{supp}(a^{in}) \subset B(0, R)$

$$\begin{aligned} \|(1 - \chi(x, -i\epsilon\partial_x))\psi_\epsilon^{in}\|_{L^2(\mathbf{R}^N)} &= \|\chi_1(1 - \chi_2(-i\epsilon\partial_x))\psi_\epsilon^{in}\|_{L^2(\mathbf{R}^N)} \\ &\leq \|(1 - \chi_2(-i\epsilon\partial_x))\psi_\epsilon^{in}\|_{L^2(\mathbf{R}^N)} \\ &= (2\pi)^{-N} \|(1 - \chi_2(\epsilon\xi))\widehat{\psi_\epsilon^{in}}\|_{L^2(\mathbf{R}^N)} \\ &\leq (2\pi)^{-N} \|\mathbf{1}_{[Q/\epsilon, \infty)}(|\xi|)\widehat{\psi_\epsilon^{in}}\|_{L^2(\mathbf{R}^N)} \end{aligned}$$

Since

$$\widehat{\psi_\epsilon^{in}}(\zeta/\epsilon) = \int_{\mathbf{R}^N} e^{-i(\zeta \cdot x - S^{in}(x))/\epsilon} a^{in}(x) dx$$

we conclude from estimate (7.7.1') in [8] that

$$|\widehat{\psi_\epsilon^{in}}(\zeta/\epsilon)| \leq \frac{C\|a^{in}\|_{W^{m,\infty}(\mathbf{R}^N)}}{(|\zeta| - \|\nabla S^{in}\|_{L^\infty(B(0,R))})^m} \epsilon^m$$

provided that $\text{supp}(a^{in}) \subset B(0, R)$ and $|\zeta| > 1 + \|\nabla S^{in}\|_{L^\infty(B(0,R))}$. Therefore

$$(26) \quad \begin{aligned} \|\psi_\epsilon(t, \cdot) - G_{\epsilon,n}(t)\chi(x, -i\epsilon\partial_x)\psi_\epsilon^{in}\|_{L^2(\mathbf{R}^N)} &\leq C_{T,Q}\|a^{in}\|_{L^2(\mathbf{R}^N)}\epsilon^{n-2N} \\ &\quad + C\|(1 + |\zeta|)^{-m}\|_{L^2(\mathbf{R}^N)}\|a^{in}\|_{W^{m,\infty}(\mathbf{R}^N)}\epsilon^m \end{aligned}$$

for all $m > N/2$.

Next we analyze the term

$$(27) \quad \begin{aligned} &G_{\epsilon,n}(t)\chi(x, -i\epsilon\partial_x)\psi_\epsilon^{in}(x) \\ &= \iiint\!\!\!\int A_n(t, y, \eta, \epsilon)a^{in}(z)\chi_2(\zeta)e^{i(\varphi(t,x,y,\eta) + \zeta \cdot (y-z) + S^{in}(z))/\epsilon} \frac{dz d\zeta dy d\eta}{(2\pi\epsilon)^{2N}} \end{aligned}$$

with the stationary phase method.

Choose φ of the form¹

$$\varphi(t, x, y, \eta) = S(t, y, \eta) + (x - X_t(y, \eta)) \cdot \Xi_t(y, \eta) + iB : (x - X_t(y, \eta))^{\otimes 2}$$

where the matrix $B = B^T > 0$ is constant (see formula (2.7) [12] and the following Remark 2.1). Critical points of the phase in the oscillating integral (27) are defined by the system of equations²

$$\begin{cases} -\zeta + DS^{in}(z) = 0, \\ y - z = 0, \\ \partial_y S(t, x, y) - D_y X_t(y, \eta)^T \cdot \Xi_t(y, \eta) + D_y \Xi_t(y, \eta)^T \cdot (x - X_t(y, \eta)) \\ - iB : (x - X_t(y, \eta)) \otimes D_y X_t(y, \eta) + \zeta = 0, \\ \partial_\eta S(t, x, y) - D_\eta X_t(y, \eta)^T \cdot \Xi_t(y, \eta) + D_\eta \Xi_t(y, \eta)^T \cdot (x - X_t(y, \eta)) \\ - iB : (x - X_t(y, \eta)) \otimes D_\eta X_t(y, \eta) = 0. \end{cases}$$

At this point, we recall formulas (3.1-2) from [12]

$$(28) \quad \begin{aligned} \partial_y S(t, y, \eta) &= D_y X_t(y, \eta)^T \cdot \Xi_t(y, \eta) - \eta, \\ \partial_\eta S(t, y, \eta) &= D_\eta X_t(y, \eta)^T \cdot \Xi_t(y, \eta), \end{aligned}$$

together with the following definitions

$$\begin{aligned} Y(t, y, \eta) &:= D_y \Xi_t(y, \eta) - iB D_y X_t(y, \eta), \\ Z(t, y, \eta) &:= D_\eta \Xi_t(y, \eta) - iB D_\eta X_t(y, \eta). \end{aligned}$$

Thus the critical points of the phase in (27) are given by

$$\begin{cases} \zeta = DS^{in}(z), \\ y = z, \\ (x - X_t(y, \eta))^T Y(t, y, \eta) + \zeta = \eta, \\ (x - X_t(y, \eta))^T Z(t, y, \eta) = 0. \end{cases}$$

Since the matrix Z is invertible by Lemma 4.1 of [12], we conclude that the system of equations above is equivalent to

$$\begin{cases} \zeta = DS^{in}(z), \\ y = z, \\ \zeta = \eta, \\ x = X_t(y, \eta). \end{cases}$$

In other words,

$$F_t(z) = x, \quad y = z, \quad \zeta = \eta = DS^{in}(z).$$

¹For each $v \in \mathbf{R}^N$, the tensor $v \otimes v$ is identified with the matrix with entries $v_i v_j$, where v_i is the i th component of the vector v in the canonical basis of \mathbf{R}^N . For $A, B \in M_N(\mathbf{R})$, the notation $A : B$ designates $\text{trace}(A^T B)$.

²If $f \in C^1(\mathbf{R}^N; \mathbf{R}^N)$, and if $f_i(x)$ designates the i th component of $f(x)$ in the canonical basis of \mathbf{R}^N , the notation $Df(x)$ designates the Jacobian matrix of f at the point x , i.e. the matrix whose entry at the i th row and the j th column is the partial derivative $\partial_{x_j} f_i(x)$.

Assuming that $(t, x) \notin C$, we apply Proposition 2.1 and conclude that the set of critical points of the phase in (27) is of the form

$$\begin{cases} y = z = y_j(t, x), \\ \zeta = \eta = DS^{in}(y_j(t, x)), \end{cases} \quad j = 1, \dots, \mathcal{N}(t, x).$$

At this point, we apply the stationary phase method (Theorem 7.7.5 in [8]). First we need to compute the Hessian of the phase in (27) at its critical points. One finds

$$H_j(t, x) := \begin{pmatrix} D^2 S^{in} & -I & 0 & 0 \\ -I & 0 & +I & 0 \\ 0 & +I & -Y^T D_y X_t & -Y^T D_\eta X_t - I \\ 0 & 0 & -Z^T D_y X_t & -Z^T D_\eta X_t \end{pmatrix}_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}}$$

and it remains to compute $\det(H_j(t, x))$. Adding the first row of $H_j(t, x)$ to the third row, one finds that

$$\begin{aligned} \det H_j(t, x) &= \begin{vmatrix} D^2 S^{in} & -I & 0 & 0 \\ -I & 0 & +I & 0 \\ D^2 S^{in} & 0 & -Y^T D_y X_t & -Y^T D_\eta X_t - I \\ 0 & 0 & -Z^T D_y X_t & -Z^T D_\eta X_t \end{vmatrix}_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}} \\ &= \begin{vmatrix} -I & +I & 0 & 0 \\ D^2 S^{in} & -Y^T D_y X_t & -Y^T D_\eta X_t - I & 0 \\ 0 & -Z^T D_y X_t & -Z^T D_\eta X_t & 0 \end{vmatrix}_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}} \\ &= \begin{vmatrix} -I & 0 & 0 & 0 \\ D^2 S^{in} & -Y^T D_y X_t + D^2 S^{in} & -Y^T D_\eta X_t - I & 0 \\ 0 & -Z^T D_y X_t & -Z^T D_\eta X_t & 0 \end{vmatrix}_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}} \end{aligned}$$

where the last equality follows from adding the first column in the right hand side of the second equality to the second column. Eventually, one finds that

$$\det H_j(t, x) = (-1)^N \begin{vmatrix} -Y^T D_y X_t + D^2 S^{in} & -Y^T D_\eta X_t - I \\ -Z^T D_y X_t & -Z^T D_\eta X_t \end{vmatrix}_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}}$$

which is computed as follows. First

$$\begin{vmatrix} I & -(Z^{-1}Y)^T \\ 0 & I \end{vmatrix} \begin{vmatrix} -Y^T D_y X_t + D^2 S^{in} & -Y^T D_\eta X_t - I \\ -Z^T D_y X_t & -Z^T D_\eta X_t \end{vmatrix} = \begin{vmatrix} D^2 S^{in} & -I \\ -Z^T D_y X_t & -Z^T D_\eta X_t \end{vmatrix}$$

so that

$$\det H_j(t, x) = (-1)^N \begin{vmatrix} D^2 S^{in} & -I \\ -Z^T D_y X_t & -Z^T D_\eta X_t \end{vmatrix}_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}}$$

On the other hand

$$\begin{aligned} (29) \quad \begin{vmatrix} D^2 S^{in} & -I \\ -Z^T D_y X_t & -Z^T D_\eta X_t \end{vmatrix} &= (-1)^N \det(Z^T D_y X_t + Z^T D_\eta X_t D^2 S^{in}) \\ &= (-1)^N \det(Z) \det(D_y X_t + D_\eta X_t D^2 S^{in}) \end{aligned}$$

by the following elementary lemma (that is a variant of the Schur complement formula in a special case: see for instance Proposition 3.9 on pp. 40-41 in [18]).

Lemma 4.2. *Let $A, B, C, D \in M_N(\mathbf{C})$. If $AB = BA$, one has*

$$\begin{vmatrix} A & B \\ C & D \end{vmatrix} = \det(DA - CB).$$

Therefore

$$(30) \quad \begin{aligned} \det H_j(t, x) &= \det(Z) \det(D_y X_t + D_\eta X_t D^2 S^{in}) \Big|_{y=y_j(t, x), \eta=DS^{in}(y_j(t, x))} \\ &= \det(Z(y_j(t, x), DS^{in}(y_j(t, x)))) \det(DF_t(y_j(t, x))) \end{aligned}$$

where F_t is defined in (9).

Pick a nonempty closed ball $\mathbf{b} \subset \mathbf{R} \times \mathbf{R}^N \setminus C$, let $\mathcal{N}_{\mathbf{b}} = \mathcal{N}(t, x)$ for all $(t, x) \in \mathbf{b}$, and let

$$K_j = \{(y_j(t, x), \nabla S^{in}(y_j(t, x))) \mid (t, x) \in \mathbf{b}\}, \quad j = 1, \dots, \mathcal{N}_{\mathbf{b}}.$$

Assuming that \mathbf{b} is of small enough radius, $K_j \cap K_k = \emptyset$ for $j \neq k \in \{1, \dots, \mathcal{N}_{\mathbf{b}}\}$. Let $\kappa_j \in C_c^\infty(\mathbf{R}^{2N})$ for all $j = 1, \dots, \mathcal{N}_{\mathbf{b}}$, such that

$$\begin{cases} \kappa_j \geq 0 \text{ and } \kappa_j|_{K_j} = 1, & j = 1, \dots, \mathcal{N}_{\mathbf{b}}, \\ \text{while } \kappa_j \kappa_k = 0 \text{ for } j \neq k \in \{1, \dots, \mathcal{N}_{\mathbf{b}}\}, \end{cases}$$

Applying Theorem 7.7.1 in [8] shows that

$$(31) \quad \begin{aligned} \sup_{(t, x) \in \mathbf{b}} \left| \iiint A_n(t, y, \eta, \epsilon) a^{in}(z) \chi_2(\zeta) e^{i(\varphi(t, x, y, \eta) + \zeta \cdot (y - z) + S^{in}(z)) / \epsilon} \frac{dz d\zeta dy d\eta}{(2\pi\epsilon)^{2N}} \right. \\ \left. - \sum_{j=1}^{\mathcal{N}_{\mathbf{b}}} \iiint A_n(t, y, \eta, \epsilon) a^{in}(z) \chi_2(\zeta) \kappa_j(y, \eta) \kappa_j(z, \zeta) \right. \\ \left. e^{i(\varphi(t, x, y, \eta) + \zeta \cdot (y - z) + S^{in}(z)) / \epsilon} \frac{dz d\zeta dy d\eta}{(2\pi\epsilon)^{2N}} \right| = O(\epsilon) \end{aligned}$$

as $\epsilon \rightarrow 0$.

Next we set

$$\begin{aligned} I_j(t, x, \epsilon) &:= \iiint A_n(t, y, \eta, \epsilon) a^{in}(z) \chi_2(\zeta) \kappa_j(y, \eta) \kappa_j(z, \zeta) \\ &\quad \times e^{i(\varphi(t, x, y, \eta) + \zeta \cdot (y - z) + S^{in}(z)) / \epsilon} \frac{dz d\zeta dy d\eta}{(2\pi\epsilon)^{2N}} \end{aligned}$$

for $j = 1, \dots, \mathcal{N}_{\mathbf{b}}$. By Theorem 7.7.5 in [8], we conclude that

$$(32) \quad \begin{aligned} \sup_{(t, x) \in \mathbf{b}} |I_j(t, x, \epsilon) - A_0(t, y_j(t, x), \nabla S^{in}(y_j(t, x)), 0) a^{in}(y_j(t, x)) \chi_2(\nabla S^{in}(y_j(t, x))) \\ \times e^{i(\varphi(t, x, y_j(t, x), \nabla S^{in}(y_j(t, x))) + S^{in}(y_j(t, x))) / \epsilon} (\det H_j(t, x))^{-1/2}| = O(\epsilon) \end{aligned}$$

as $\epsilon \rightarrow 0$. Our choice of χ_2 and φ implies that $\chi_2(\nabla S^{in}(y_j(t, x))) = 1$ and

$$\varphi(t, x, y_j(t, x), \nabla S^{in}(y_j(t, x))) = S(t, y_j(t, x), \nabla S^{in}(y_j(t, x)))$$

so that

$$(\varphi(t, x, y_j(t, x), \nabla S^{in}(y_j(t, x))) + S^{in}(y_j(t, x))) = S_j(t, x).$$

By formula (2.13) in [12]

$$(33) \quad A_0(t, y_j(t, x), \nabla S^{in}(y_j(t, x)), 0) = \sqrt[cont]{\det(Z(t, y_j(t, x), DS^{in}(y_j(t, x))))},$$

where the notation ${}^{\text{con}}\sqrt{z}$ designates the analytic continuation of the square-root along the path $s \mapsto \det(Z(s, y_j(t, x), DS^{in}(y_j(t, x)))$. This analytic continuation is uniquely defined since $\det(Z(s, y_j(t, x), DS^{in}(y_j(t, x))) \neq 0$ for all s : see for instance section 1.3 in chapter 8 of [1]. According to (33), (30), Lemma 5.1 and formula (5.15) in [12], we have

$$(34) \quad \begin{aligned} & A_0(t, y_j(t, x), \nabla S^{in}(y_j(t, x)), 0)(\det H_j(t, x))^{-1/2} \\ &= |\det(DF_t(y_j(t, x)))|^{-1/2} e^{i\pi\nu_j(t, x)/2} = J(y_j(t, x))^{-1/2} e^{-i\pi\mathcal{M}_j(t, x)/2}, \end{aligned}$$

where $\mathcal{M}_j(t, x)$ is an integer.

Putting together (26)-(31)-(32)-(34) concludes the proof of Theorem 4.1. \square

Proof of Lemma 4.2. If A is nonsingular and $AB = BA$, one has

$$\begin{aligned} \begin{vmatrix} A & B \\ C & D \end{vmatrix} &= (-1)^N \begin{vmatrix} A & B \\ C & D \end{vmatrix} \begin{vmatrix} A^{-1} & B \\ 0 & -A \end{vmatrix} = (-1)^N \begin{vmatrix} I & 0 \\ CA^{-1} & CB - DA \end{vmatrix} \\ &= (-1)^N \det(CB - DA) = \det(DA - CB). \end{aligned}$$

Since both sides of the identity above are continuous functions of A and the set of nonsingular matrices $GL_N(\mathbf{C})$ is dense in $M_N(\mathbf{C})$, this identity holds for all $A \in M_n(\mathbf{C})$ such that $AB = BA$. \square

In the case of a linear phase function S^{in} treated in section 5 of [12], one has $\nabla^2 S^{in} \equiv 0$. Therefore, the determinant in (29) reduces to

$$\begin{vmatrix} 0 & -I \\ -Z^T D_y X_t & -Z^T D_\eta X_t \end{vmatrix}$$

which can be explicitly computed without difficulty, since this determinant is block-wise triangular. In particular, the simpler situation considered in section 5 of [12] does not require using Lemma 4.2.

5. MASLOV-INDEX MIT MENSCHLICHEM ANTLITZ

The purpose of this section is to explain, in the simplest possible manner, how the integers $\mathcal{M}_j(t, x)$ appearing in formulas (15) and (24) are related to the Maslov index.

5.1. Generalities on the Maslov index. There are several notions of Maslov index in the literature. The original definition can be found in Maslov's treatise [15] or [2]. Closely related indices have been subsequently defined by Leray (in §2 of [13]) and Hörmander (in chapter XXI of [10]) — see also chapters I.7 and IV.3 in [7], and especially [19] for a lucid presentation of all these notions of Maslov index and how they are related. The present section recalls some material from Arnold's short and precise presentation [2] of the subject.

The phase space $\mathbf{R}_x^N \times \mathbf{R}_\xi^N$ is endowed with the standard symplectic 2-form $\sigma := d\xi_1 \wedge dx_1 + \dots + d\xi_N \wedge dx_N$. A linear subspace λ of $\mathbf{R}_x^N \times \mathbf{R}_\xi^N$ is called Lagrangian if and only if $\dim \lambda = N$ and $\sigma(u, v) = 0$ for all $u, v \in \lambda$. An example of Lagrangian subspace of $\mathbf{R}_x^N \times \mathbf{R}_\xi^N$ is $T_0^* := \{0\} \times \mathbf{R}^N$. The Lagrangian Grassmanian $\Lambda(N)$ is the set of Lagrangian subspaces of $\mathbf{R}_x^N \times \mathbf{R}_\xi^N$. For each $k = 0, \dots, N$ define $\Lambda^k(N) := \{\lambda \in \Lambda(N) \mid \dim(\lambda \cap T_0^*) = k\}$. A linear subspace λ of $\mathbf{R}_x^N \times \mathbf{R}_\xi^N$ belongs to $\Lambda^0(N)$ iff λ is defined by an equation of the form $\xi = Ax$ with $A = A^T \in M_N(\mathbf{R})$. The Lagrangian Grassmanian $\Lambda(N)$ is a C^∞ manifold of dimension $\frac{1}{2}N(N+1)$

(Corollary 3.1.4 in [2]), and $\Lambda^k(N)$ is a submanifold of codimension $\frac{1}{2}k(k+1)$ in $\Lambda(N)$ (Lemma 3.2.1 in [2]). An important subset of the Lagrangian Grassmanian is the *Maslov cycle* $\mathcal{M} := \Lambda(N) \setminus \Lambda^0(N) = \Lambda^1(N) \cup \dots \cup \Lambda^N(N)$; it has codimension 1 in $\Lambda(N)$ (see section 3.2, especially Corollary 3.2.2 in [2] for a proof that the homological boundary of \mathcal{M} is 0).

Let us define an orientation on \mathcal{M} . For each $\theta \in \mathbf{R}$ and each $(x, \xi) \in \mathbf{R}_x^N \times \mathbf{R}_\xi^N$, define $R[\theta](x, \xi) = (x \cos \theta + \xi \sin \theta, \xi \cos \theta - x \sin \theta)$. Since $R[\theta]$ is the Hamiltonian flow of $H(x, \xi) := \frac{1}{2}(|x|^2 + |\xi|^2)$, it defines a symplectomorphism of the phase space $\mathbf{R}_x^N \times \mathbf{R}_\xi^N$. In particular $R[\theta]\lambda \in \Lambda(N)$ for each $\lambda \in \Lambda(N)$ and each $\theta \in \mathbf{R}$. The Maslov cycle is oriented by the following prescription: let $\lambda \in \Lambda^1(N)$; then the path $\theta \mapsto R[\theta]\lambda$ crosses \mathcal{M} exactly once for θ near 0, at λ for $\theta = 0$, and does so from the negative side of \mathcal{M} to the positive side of \mathcal{M} , as θ increases in $(-\eta, \eta)$ for $\eta > 0$ small enough (section 3.5, especially Lemmas 3.5.1, 3.5.2 and 3.5.3 in [2]).

Let now $[0, t] \ni s \mapsto \lambda(s) \in \Lambda(N)$ be a C^1 path such that $\lambda(0), \lambda(t) \in \Lambda^0(N)$ and such that $\{\lambda(s) \mid 0 < s < t\} \cap \mathcal{M} \subset \Lambda^1(N)$ with transverse intersection. The Maslov index $\mu(\lambda)$ of the path λ is its intersection index with the Maslov cycle \mathcal{M} oriented as above. In other words

$$\mu(\lambda) = \sum_{\lambda(s) \in \mathcal{M}} \text{sign}(s),$$

where $\text{sign}(s) = +1$ if $\lambda(s+t)$ crosses \mathcal{M} from the negative to the positive side of \mathcal{M} as t increases near 0, and $\text{sign}(s) = -1$ if $\lambda(s+t)$ crosses \mathcal{M} from the positive to the negative side of \mathcal{M} as t increases near 0 (see section 2.2 and Definition 3.6.1 in [2]).

There exists an alternate definition of the Maslov index for closed paths in $\Lambda(N)$. Identifying $\mathbf{R}_x^N \times \mathbf{R}_\xi^N$ with $\mathbf{C}^N = \mathbf{R}_\xi^N + i\mathbf{R}_x^N$, we recall that the unitary group $U(N)$ acts transitively on $\Lambda(N)$ (Lemma 1.2 in [2]). Thus, for each $\lambda \in \Lambda(N)$, there exists $u \in U(N)$ such that $\lambda = uT_0^*$; besides, if $u, u' \in U(N)$ and $uT_0^* = u'T_0^*$, then $uu^T = u'(u')^T$. In other words, one can identify λ with uu^T where u is any element of $U(N)$ such that $uT_0^* = \lambda$. This defines a map $\det^2 : \Lambda(N) \rightarrow \mathbf{S}^1$ by $\det^2(\lambda) := \det(u)^2$ where $uT_0^* = \lambda$. Let now $\mathbf{S}^1 \ni s \mapsto \lambda(s) \in \Lambda(N)$ be a closed continuous path; the Maslov index $\mu(\lambda)$ of λ is defined as the winding number of the composed map $\det^2 \circ \lambda : \mathbf{S}^1 \rightarrow \mathbf{S}^1$, i.e.

$$\mu(\lambda) = \text{degree}(\det^2 \circ \lambda).$$

(See section 1.5 in [2]). If λ is a C^1 closed path on $\Lambda(N)$ intersecting \mathcal{M} transversally on $\Lambda^1(N)$, both definitions of the Maslov index coincide (Theorem 1.5 in [2]).

5.2. The Maslov index and Hamiltonians of the form $\frac{1}{2}|\xi|^2 + V(x)$. The orientation of the Maslov cycle \mathcal{M} is obviously crucial in the definition of the Maslov index recalled above. For that reason, computing the Maslov index of a path is in general a rather complicated task. However, when the path is defined by the linearized Hamiltonian flow of a Hamiltonian that is convex in the momentum variable, this computation is considerably simplified. Indeed, in that case, the orientation of the Maslov cycle plays no role as such a path always crosses the Maslov cycle in the same direction. Therefore, computing the Maslov index of such a path reduces to counting how many times it intersects the Maslov cycle. In other words, the Maslov index reduces to the more classical notion of Morse index (see for instance [17] §15, especially Theorem 15.1) in this case. This observation can

be found in [15], first without proof on p. 151, and as the result of a rather lengthy argument on p. 297. See also Theorem 5.2, given without proof in [2]. The lemma below gives a short proof of this fact in the special case of a Hamiltonian of the form $\frac{1}{2}|\xi|^2 + V(x)$, which is all that we need in the context of the classical limit of the Schrödinger equation.

Let $W \in C_b(\mathbf{R}; M_N(\mathbf{R}))$ such that $W(t) = W(t)^T$ for each $t \in \mathbf{R}$, and let $t \mapsto S(t, t_0) \in M_{2N}(\mathbf{R})$ be the solution of the Cauchy problem

$$\frac{d}{dt}S(t, t_0) = \begin{pmatrix} 0 & I \\ W(t) & 0 \end{pmatrix} S(t, t_0), \quad S(t_0, t_0) = I.$$

(The linearized Hamiltonian system defined by the Hamiltonian $\frac{1}{2}|\xi|^2 + V(x)$ and the symplectic form $d\xi_1 \wedge dx_1 + \dots + d\xi_N \wedge dx_N$ is exactly of this form, with $W = -\nabla^2 V$.) The matrix $S(t, t_0)$ is symplectic because

$$\begin{pmatrix} 0 & -I \\ I & 0 \end{pmatrix} \begin{pmatrix} 0 & W(t) \\ I & 0 \end{pmatrix} + \begin{pmatrix} 0 & I \\ W(t) & 0 \end{pmatrix} \begin{pmatrix} 0 & -I \\ I & 0 \end{pmatrix} = 0.$$

In particular $S(t, t_0)\lambda \in \Lambda(N)$ whenever $\lambda \in \Lambda(N)$.

Lemma 5.1. *Let $\lambda_0 \in \Lambda^0(N)$ and set $\lambda(t) = S(t, t_0)\lambda_0$. If $\lambda(t_1) = S(t_1, t_0)\lambda_0 \in \Lambda^1(N)$, the path $t \mapsto \lambda(t)$ intersects the Maslov cycle \mathcal{M} transversally at $\lambda(t_1)$ from the negative side to the positive side of the Maslov cycle as t increases near t_1 .*

Proof. Choose a system of orthonormal coordinates p_1, \dots, p_N in \mathbf{R}_ξ^N such that the line $\lambda(t_1) \cap T_0^*$ is transverse to the hyperplane of equation $p_1 = 0$ in T_0^* . Choose orthonormal coordinates q_1, \dots, q_N in \mathbf{R}_x^N that are conjugate to p_1, \dots, p_N , i.e. such that the symplectic form $\sigma = dp_1 \wedge dq_1 + \dots + dp_N \wedge dq_N$. The change of coordinates takes the form

$$\begin{pmatrix} R & 0 \\ 0 & R \end{pmatrix}$$

where $R \in O_N(\mathbf{R})$. In these new coordinates, the differential equation defining $S(t, t_0)$ keeps the same form, up to replacing $W(t)$ with $RW(t)R^T$. For simplicity, we keep the same notation for $S(t, t_0)$ and $W(t)$ in these new variables.

Define

$$\hat{p}_1 = q_1, \quad \hat{q}_1 = -p_1, \quad \hat{p}_j = p_j \text{ and } \hat{q}_j = q_j \quad \text{for } j = 2, \dots, N.$$

In these coordinates, any Lagrangian space λ' such that $T_0^* \cap \lambda'$ is a line transverse to the hyperplane of equation $p_1 = 0$ in T_0^* is defined by an equation of the form $\hat{p} = L'\hat{q}$ with $L' = (L')^T \in M_N(\mathbf{R})$ such that $L'_{11} = 0$.

Denote $I_1 := \text{diag}(1, 0, \dots, 0) \in M_N(\mathbf{R})$ and $I' := I - I_1$, and set

$$J := \begin{pmatrix} I' & I_1 \\ -I_1 & I' \end{pmatrix}$$

so that $J(q, p)^T = (\hat{q}, \hat{p})^T$. Straightforward computations show that

$$JS(t_1 + \tau, t_1)J^T = \begin{pmatrix} I - \tau I_1 W(t_1) I' & \tau I' + \tau I_1 W(t_1) I_1 \\ -\tau I' W(t_1) I' - \tau I_1 & I + \tau I' W(t_1) I_1 \end{pmatrix} + O(\tau)^2.$$

Let $L = L^T \in M_N(\mathbf{R})$ be such that the Lagrangian subspace $\lambda(t_1)$ has equation $\hat{p} = L\hat{q}$. For $|\tau| \ll 1$, the Lagrangian subspace $\lambda(t_1 + \tau)$ has equation $\hat{p} = L(\tau)\hat{q}$,

where

$$\begin{aligned} L(\tau) = & (-\tau I'W(t_1)I' - \tau I_1 + (I + \tau I'W(t_1)I_1)L) \\ & \times (I - \tau I_1W(t_1)I' + (\tau I' - \tau I_1W(t_1)I_1)L)^{-1} + O(\tau^2) \end{aligned}$$

so that, in the limit as $\tau \rightarrow 0$, one has

$$\begin{aligned} \frac{1}{\tau}(L(\tau) - L) \rightarrow & -I'W(t_1)I' - I_1 + I'W(t_1)I_1L \\ & + LI_1W(t_1)I' - LI'L - LI_1W(t_1)I_1L. \end{aligned}$$

Assume that $\lambda(t_1) \in \Lambda^1(N)$ and that $T_0^* \cap \lambda(t_1)$ is transverse to the hyperplane of equation $\hat{p}_1 = 0$ on T_0^* , so that $L_{11} = 0$. Let e_1 be the first vector in the canonical basis of \mathbf{R}^N ; then, in the limit as $\tau \rightarrow 0$, one has

$$\frac{1}{\tau}(e_1|(L(\tau) - L)e_1) = \frac{1}{\tau}(e_1|L(\tau)e_1) \rightarrow -1 - |I'Le_1|^2 \leq -1.$$

In the special case $W = -I$, one has $S(t, t_0) = R[t - t_0]$. Therefore, near $\lambda(t_1)$, the computation above shows that the positive side of $\Lambda^1(N)$ consists of Lagrangian subspaces λ' of equation $\hat{p} = L'\hat{q}$ with $L'_{11} < 0$.

Besides, the computation above also shows that, for all W , the path $t \mapsto S(t, t_1)\lambda(t_1)$ crosses $\Lambda^1(N)$ in the same direction as for $W = -I$, i.e. from the negative side to the positive side as t increases near t_1 . \square

5.3. The free case. Let $A = A^T \in GL_N(\mathbf{R})$, and let $\lambda(0)$ be the Lagrangian subspace of equation $\xi = Ax$ in $\mathbf{R}_x^N \times \mathbf{R}_\xi^N$. Let $\Phi_t(x, \xi) = (x + t\xi, \xi)$ be the free flow defined on $\mathbf{R}_x^N \times \mathbf{R}_\xi^N$ for each $t \in \mathbf{R}$, which is the Hamiltonian flow of $H(x, \xi) = \frac{1}{2}|\xi|^2$. Set $\lambda(s) := \Phi_s(\lambda(0)) \in \Lambda(N)$ for each $s \in \mathbf{R}$.

First assume that A has N distinct eigenvalues $\alpha_1 > \dots > \alpha_N$, and let $t > 0$ be such that $I + tA$ is invertible. Consider the path $\lambda_t : [0, t] \ni s \mapsto \lambda(s) \in \Lambda(N)$. Obviously $\lambda_t(s) := \{(x + sAx, Ax) \mid x \in \mathbf{R}^N\}$. Hence $\lambda_t(s) \in \Lambda^0(N)$ if $I + sA$ is invertible, and has equation $\xi = A(I + sA)^{-1}x$.

If $I + sA$ is not invertible, then $\text{Ker}(I + sA)$ has dimension 1 since A has simple eigenvalues, and therefore $\lambda_t(s) \cap T_0^* = \{0\} \times \text{Ker}(I + sA)$ has dimension 1. Thus the path λ_t can only intersect the Maslov cycle \mathcal{M} on its regular part $\Lambda^1(N)$. By Lemma 5.1, it always does so in the same direction, from the negative to the positive side of \mathcal{M} . Hence the Maslov index of the path λ_t is

$$\mu(\lambda_t) = \#\{\alpha_j \mid 0 < -1/\alpha_j < t\} = \#\{\alpha_j \mid 1 + t\alpha_j < 0\}.$$

Next we treat the general case, where A may have multiple eigenvalues, still denoted $\alpha_1 \geq \dots \geq \alpha_N$ and counted with their multiplicities, and compute the Maslov index $\mu(\lambda_t)$ of the path λ_t defined above.

Pick $A' = (A')^T$ near A with distinct eigenvalues and such that $A'A = AA'$. Set $A(\tau) = (1 - \tau)A + \tau A'$. Assume that $I + tA$ is invertible; by choosing A' sufficiently close to A , one can assume that $I + tA(\tau)$ is invertible for each $\tau \in [0, 1]$. For each $B = B^T \in M_N(\mathbf{R})$, denote by $\lambda[B]$ the Lagrangian subspace of equation $\xi = Bx$. Consider now the family indexed by $\tau \in [0, 1]$ of closed paths γ_τ defined as follows

$$\begin{aligned} \gamma_\tau(s) &= \Phi_{ts}\lambda[A], & \text{for } 0 \leq s \leq 1, \\ \gamma_\tau(s) &= \Phi_t\lambda[A(\tau(s-1))], & \text{for } 1 < s < 2, \\ \gamma_\tau(s) &= \Phi_{t(3-s)}\lambda[A(\tau)], & \text{for } 2 \leq s \leq 3, \\ \gamma_\tau(s) &= \lambda[A(\tau(4-s))], & \text{for } 3 < s < 4. \end{aligned}$$

Letting $\tau \rightarrow 0$ shows that γ_1 is homotopic to $\lambda_t - \lambda_t$ (i.e. the path λ_t followed by its opposite). By the homotopy invariance of the degree, using the definition of Maslov's index for closed continuous paths shows that $\mu(\gamma_\tau) = 0$. Since $I + tA(\tau)$ is invertible for $\tau \in [0, 1]$, the arcs of γ_1 corresponding to $s \in [1, 2]$ and $s \in [3, 4]$ never cross \mathcal{M} . Therefore

$$\mu(\lambda_t) = -\mu(\gamma_1|_{[2,3]}),$$

from which we conclude that

$$\mu(\lambda_t) = \#\{\alpha_j \mid 1 + t\alpha_j < 0\}.$$

5.4. The non free case. Denote by $\Phi_t = (X_t, \Xi_t)$ the Hamiltonian flow of

$$H(x, \xi) := \frac{1}{2}|\xi|^2 + V(x),$$

let $(t, x) \notin C$, and let $\mathcal{L}_0 := \{(y, \nabla S^{in}(y)) \text{ s.t. } y \in \mathbf{R}^N\}$. Denote $\mathcal{L}_s := \Phi_s(\mathcal{L}_0)$ for each $s \in \mathbf{R}$. Set $\gamma(s) := \Phi_s(y_j(t, x), \nabla S^{in}(y_j(t, x)))$ for each $s \in \mathbf{R}$ (where y_j has been defined in Proposition 2.1, and $\lambda_0 := T_{\gamma(0)}\mathcal{L}_0$. Let $\lambda_s := D\Phi_s(\gamma(0))\lambda_0$ for each $s \in \mathbf{R}$; observe that $\gamma(s + s') = \Phi_{s'}(\gamma(s))$ and that $D\Phi_{s'}(\gamma(s))\lambda(s) := \lambda(s + s')$ by the chain rule. With this notation, the point $X_s(y_j(t, x), \nabla S^{in}(y_j(t, x)))$ belongs to C_s if and only if $\lambda(s) \in \bar{\Lambda}^1(N)$.

Assume that $\lambda(s) \in \Lambda^1(N)$ whenever $\gamma(s) \in C_s$ for $0 < s < t$.

In that case, by Lemma 5.1, the path $s \mapsto \lambda(s)$ always crosses the Maslov cycle transversally from the negative to the positive side as $\gamma(s) \in C_s$. Therefore, the Maslov index of the path $[0, t] \ni s \mapsto \lambda(s) \in \Lambda(N)$ is in this case

$$\mu((\lambda(s))_{0 \leq s \leq t}) = \#\{s \in [0, t] \text{ s.t. } X_s(y_j(t, x), \nabla S^{in}(y_j(t, x))) \in C_s\}.$$

The set $\{s \in [0, t] \text{ s.t. } X_s(y_j(t, x), \nabla S^{in}(y_j(t, x))) \in C_s\}$ is therefore finite and henceforth denoted by $0 < s_1 < s_2 < \dots < s_n < t$. For $k = 1, \dots, n$, there exists $\phi_k \neq 0$ such that $\text{Ker}(DF_{s_k}(y_j(t, x))) = \mathbf{R}\phi_k$.

Consider on the other hand the matrix

$$M(s) := \begin{pmatrix} -Y_s^T D_y X_s + D^2 S^{in} & -Y_s^T D_\eta X_s - I \\ -Z_s^T D_y X_s & -Z_s^T D_\eta X_s \end{pmatrix} \Big|_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}}$$

where

$$Y_s = D_y \Xi_s - iBD_y X_s, \quad \text{and } Z_s = D_\eta \Xi_s - iBD_\eta X_s.$$

Denote

$$M_1(s) := \begin{pmatrix} -D_y \Xi_s^T D_y X_s + D^2 S^{in} & -D_y \Xi_s^T D_\eta X_s - I \\ -D_\eta \Xi_s^T D_y X_s & -D_\eta \Xi_s^T D_\eta X_s \end{pmatrix} \Big|_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}}$$

and

$$M_2(s) := \begin{pmatrix} D_y X_s^T \\ D_\eta X_s^T \end{pmatrix} B(D_y X_s \ D_\eta X_s) \Big|_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}}$$

so that

$$M_1(s) = M_1(s)^T = \Re(M(s)), \quad M_2(s) = M_2(s)^T = \Im(M(s, \theta)) \geq 0.$$

Since $\det(M(s)) = (-1)^N \det(Z_s) \det(DF_s(y_j(t, x)))$ and Z_s is invertible for each $s \in [0, t]$, one has $\det(M(s)) = 0$ at $s = s_1 < s_2 < \dots < s_n$ only in $[0, t]$. Besides,

$$\begin{aligned} M(s_k)(\phi \oplus \psi) = 0 &\Rightarrow \psi = D^2 S^{in}(y_j(t, x))\phi \text{ and } DF_{s_k}(y_j(t, x))\phi = 0 \\ &\Rightarrow M_1(s_k)(\phi \oplus \psi) = M_2(s_k)(\phi \oplus \psi) = 0, \end{aligned}$$

the converse being obvious. Hence

$$\text{Ker}(M(s_k)) = \mathbf{C}(\phi_k \oplus D^2 S^{in}(y_j(t, x))\phi_k),$$

while

$$\text{Ker}(M_1(s_k)) \cap \text{Ker}(M_2(s_k)) = \mathbf{R}(\phi_k \oplus D^2 S^{in}(y_j(t, x))\phi_k).$$

Denote $V_k := \mathbf{R}(\phi_k \oplus D^2 S^{in}(y_j(t, x))\phi_k)$; since $M_j(s_k)$ is a real symmetric matrix and $M_j(s_k)(V_k) \subset V_k$ for $j = 1, 2$, one has $M_j(s_k)(V_k^\perp) \subset V_k^\perp$. Consider now the linear space $W_k := V_k^\perp \oplus iV_k^\perp$; one has $M(s_k)W_k \subset W_k$ and

$$\mathbf{C}^{2N} = \text{Ker}(M(s_k, \theta)) \oplus W_k.$$

In particular $M(s_k)|_{W_k}$ is invertible on W_k . Therefore, computing the characteristic polynomial of $M(s_k)$ in this decomposition of \mathbf{C}^{2N} , we find that

$$\det(\lambda I_{2N} - M(s_k)) = \lambda \det(\lambda I_{2N-1} - M(s_k)|_{W_k}),$$

so that $\lambda = 0$ is a simple root of the characteristic polynomial of $M(s_k)$.

By the implicit function theorem, there exists a local C^1 function $s \mapsto \lambda_k(s)$ defined near s_k and such that

$$\lambda_k(s_k) = 0, \quad \text{and } \det(\lambda_k(s)I_{2N} - M(s)) = 0 \text{ for all } s \text{ near } s_k.$$

Moreover, there exist a C^1 vector field $s \mapsto \psi_k(s) \neq 0$ defined near s_k such that

$$M(s)\psi_k(s) = \lambda_k(s)\psi_k(s), \quad \psi_k(s_k) = \phi_k \oplus D^2 S^{in}(y_j(t, x))\phi_k.$$

Differentiating in s at $s = s_k$, one finds

$$\dot{M}(s_k)\psi_k(s_k) + M(s_k)\dot{\psi}_k(s_k) = \dot{\lambda}_k(s_k)\psi_k(s_k)$$

and observing that $M(s) = M(s)^T$, one concludes that

$$\psi_k(s_k)^T \dot{M}(s_k)\psi_k = \dot{\lambda}_k(s_k)\psi_k(s_k)^T \psi_k(s_k),$$

so that

$$\dot{\lambda}_k(s_k) = \frac{\psi_k(s_k)^T \dot{M}(s_k)\psi_k(s_k)}{|\psi_k(s_k)|^2}.$$

Notice that $\psi_k(s_k)^T \psi_k(s_k) = |\psi_k(s_k)|^2 > 0$ since $\psi_k(s_k) = \phi_k \oplus D^2 S^{in}(y_j(t, x))\phi_k$ belongs to $\mathbf{R}^{2N} \setminus \{0\}$.

Observe that

$$\partial_s D_y X_s = D_y \Xi_s, \quad \text{and } \partial_s D_\eta X_s = D_\eta \Xi_s,$$

so that

$$\begin{aligned} \dot{M}(s_k) &= - \begin{pmatrix} \partial_s Y_{s_k}^T \\ \partial_s Z_{s_k}^T \end{pmatrix} (D_y X_{s_k} \quad D_\eta X_{s_k}) \Big|_{\substack{y=y_j(t,x) \\ \eta=D^2 S^{in}(y_j(t,x))}} \\ &\quad - \begin{pmatrix} Y_{s_k}^T \\ Z_{s_k}^T \end{pmatrix} (D_y \Xi_{s_k} \quad D_\eta \Xi_{s_k}) \Big|_{\substack{y=y_j(t,x) \\ \eta=D^2 S^{in}(y_j(t,x))}} \end{aligned}$$

and

$$\dot{M}(s_k)\psi_k(s_k) = - \begin{pmatrix} Y_{s_k}^T \\ Z_{s_k}^T \end{pmatrix} (D_y \Xi_{s_k} \quad D_\eta \Xi_{s_k}) \Big|_{\substack{y=y_j(t,x) \\ \eta=D^2 S^{in}(y_j(t,x))}} \psi_k(s_k).$$

Thus

$$\begin{aligned} \psi_k(s_k)^T \dot{M}(s_k) \psi_k(s_k) &= i \psi_k(s_k)^T \begin{pmatrix} D_y X_{s_k}^T \\ D_\eta X_{s_k}^T \end{pmatrix} B(D_y \Xi_{s_k} \ D_\eta \Xi_{s_k}) \Big|_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}} \psi_k(s_k) \\ &\quad - \psi_k(s_k)^T \begin{pmatrix} D_y \Xi_{s_k}^T \\ D_\eta \Xi_{s_k}^T \end{pmatrix} (D_y \Xi_{s_k} \ D_\eta \Xi_{s_k}) \Big|_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}} \psi_k(s_k). \end{aligned}$$

Since $\psi_k(s_k) = \phi_k \oplus D^2 S^{in}(y_j(t, x)) \phi_k$ with $DF_{s_k}(y_j(t, x)) \phi_k = 0$, the first term on the right hand side is

$$\begin{aligned} &i \psi_k(s_k)^T \begin{pmatrix} D_y X_{s_k}^T \\ D_\eta X_{s_k}^T \end{pmatrix} B(D_y \Xi_{s_k} \ D_\eta \Xi_{s_k}) \Big|_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}} \psi_k(s_k) \\ &= i \phi_k^T DF_{s_k}(y_j(t, x))^T B(D_y \Xi_{s_k} \ D_\eta \Xi_{s_k}) \Big|_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}} \psi_k(s_k) = 0. \end{aligned}$$

Hence

$$\psi_k(s_k)^T \dot{M}(s_k) \psi_k(s_k) = -|(D_y \Xi_{s_k} \ D_\eta \Xi_{s_k}) \psi_k(s_k)|^2 \leq 0.$$

This last inequality can obviously not be an equality since

$$(D_y X_{s_k} \ D_\eta X_{s_k}) \Big|_{\substack{y=y_j(t,x) \\ \eta=DS^{in}(y_j(t,x))}} \psi_k(s_k) = DF_{s_k}(y_j(t, x)) \phi_k = 0,$$

while

$$\det \begin{pmatrix} D_y X_s & D_\eta X_s \\ D_y \Xi_s & D_\eta \Xi_s \end{pmatrix} = 1$$

and $\psi_k(s_k) \neq 0$. Therefore

$$\dot{\lambda}_k(s_k) < 0.$$

Now, the function $[0, t] \ni s \mapsto \sqrt{\det(M(s)/i)} \in \mathbf{C}$ is continuous, and its argument has jump discontinuities for $s = s_k$ for $k = 1, \dots, n$ only. Each time $s \in [0, t]$ crosses one of the values s_k , the jump in the argument of $\sqrt{\det(M(s)/i)}$ is exactly the jump in the argument of $\sqrt{\lambda_k(s)/i}$, and the previous computation shows that this jump is exactly $+\frac{\pi}{2}$. Hence

$$\mathcal{M}_j(t, x) = n = \mu((\lambda(s))_{0 \leq s \leq t}).$$

REFERENCES

- [1] L.V. Ahlfors: “Complex Analysis”, McGraw Hill, 2nd edition, New York 1966.
- [2] V.I. Arnold: *Characteristic class entering in quantization condition*, Func. Anal. Appl. **1** (1967), 1–14.
- [3] V.I. Arnold: “Geometrical methods of the theory of ordinary differential equations”, Springer-Verlag, New York 1997.
- [4] V.I. Arnold: “Mathematical Methods of Classical Mechanics”, Springer-Verlag, New York 1989.
- [5] C. Bardos, F. Golse, P. Markowich, T. Paul: *Hamiltonian Evolution of Monokinetic Measures with Rough Momentum Profile*, preprint arXiv:1207.5927.
- [6] P. Gérard, P. Markowich, N. Mauser, F. Poupaud: *Homogenization limit and Wigner Transforms*, Comm. on Pure and App. Math. **50** (1997), 323–379.
- [7] V. Guillemin, S. Sternberg: “Geometric Asymptotics”, Amer. Math. Soc., Providence 1977.
- [8] L. Hörmander: “The analysis of linear partial differential operators I”, Springer-Verlag, Berlin, Heidelberg 1983, 1990.
- [9] L. Hörmander: “The analysis of linear partial differential operators II”, Springer-Verlag, Berlin, Heidelberg 1983, 1990.
- [10] L. Hörmander: “The analysis of linear partial differential operators III”, Springer-Verlag, Berlin, Heidelberg 1985, 1994.

- [11] L. Hörmander: “The analysis of linear partial differential operators IV”, Springer-Verlag, Berlin, Heidelberg 1985, 1994.
- [12] A. Laptev, I. Sigal: *Global Fourier Integral Operators and Semiclassical Asymptotics*, Review of Math. Phys. **12** (2000), 749–766.
- [13] J. Leray: *Lagrangian Analysis and Quantum Mechanics*, The MIT Press, Cambridge Mass., 1981.
- [14] P.-L. Lions, T. Paul: *Sur les mesures de Wigner*, Rev. Mat. Iberoamericana **9** (1993), 553–618.
- [15] V.P. Maslov: *Théorie des perturbations et méthodes asymptotiques*, Dunod, Paris, 1972.
- [16] V.P. Maslov, M.V. Fedoryuk: “Semiclassical approximation in quantum mechanics”, Reidel Publishing Company, Dordrecht, 1981.
- [17] J. Milnor: “Morse Theory”, Princeton Univ. Press, Princeton NJ, 1963, 1969.
- [18] D. Serre: “Matrices”, Springer-Verlag, 2nd edition, New York 2010.
- [19] J.-M. Souriau: *Construction explicite de l’indice de Maslov*, 117–148, in “Group Theoretical Methods in Physics”, A. Janner, T. Janssen and M. Boon eds., Lecture Notes in Phys. 50, Springer-Verlag, Berlin, Heidelberg, 1976.

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